# Quantifying clumsiness in a Leggett-Garg test

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#### Invasivity of measurements in phase space

• Assuming an initial definite phase-space point  $\mathbf{r}_{in} = (x, p)$ 

• We formalize the effect of measurement as

$$\mathbf{r}_{\mathrm{in}} = (x, p) \mapsto \mathcal{M}(x, p) = (x', p') := \mathbf{r}_{\mathrm{out}} ,$$

• For initial *uncertainty* we have a PDF  $\rho(x, p)$  and

$$\rho_{\rm in}(x,p) \mapsto \rho_{\rm out}(x,p) = \int d\lambda \mu(\lambda) \rho_{\rm in}(\mathcal{M}_{\lambda}(x,p)),$$

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#### Coarse-grained dichotomic observables

Macroscopic dichotomic observables can be defined by coarse-graning

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  $P = \operatorname{sgn}(p)$ 

#### • And discrete probabilities $\Pr(Q, P)$ like

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# Macrorealism conditions

• *Macrorealism* (MR) means that  $\mathcal{M}(\mathbf{r}) = \mathbf{r}$  and in particular

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Necessary conditions follow

$$W := \sum_{Q_1, Q_3 = \pm 1} \left| \Pr(Q_1, Q_3)_{\mathcal{S}_{1,2,3}} - \Pr(Q_1, Q_3)_{\mathcal{S}_{1,3}} \right| = 0 \tag{2}$$

• On joint probabilities  $Pr(Q_i, Q_j)_{S_{(i,j)}}$  of  $S_{i,j} = \mathcal{M}_j \circ \mathcal{M}_i$ 



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A test can witness failure of MR due to

**(1)** Non-existence of a definite state  $\mathbf{r}(t)$  at  $t_1, t_2, t_3$ 

2 Invasive effect of measurement  $\mathcal{M}_i(\mathbf{r}) \neq \mathbf{r}$ 

• QM assumes both (1) and (2)

• However, also *clumsy* measurements are invasive!

• One must distinguish clumsiness from **invasivity** and quantify the former (Remember:macroscopic observables should be considered)

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# Our proposal: perform QND measurements...



# $[H_I, x_S] = 0$ $\implies \text{Quantum Non-Demolition}^1 \text{ measurement of } x_S$

#### It is suitable for macroscopic variables

[<sup>1</sup> P. Grangier, J. A. Levenson, and J.-P. Poizat, Nature 396, 537 (1998); V. B. Braginsky, Y. I. Vorontsov, and K. S. Thorne, Science 209 547 (1980)] 🗤 🔍 🔿

#### Where does MR fail?

$$p_S^{(\text{out})} = p_S^{(\text{in})} + \kappa x_M^{(\text{in})}$$

there is a **back action** on  $p_S$ 

#### Where does MR fail?

$$(\Delta x_S)^2_{(\text{out})} = \chi (\Delta x_S)^2_{(\text{in})} + (1 - \chi) (\Delta x_S)^2_{(\text{noise})}$$

and also **noise** directly in  $x_S$ 

#### Where does MR fail?

We can make an explicit distinction

 $\mathcal{M}(\mathbf{r}) = \mathcal{M}_X \circ \mathcal{M}_P \; ,$ 

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where  $\mathcal{M}_X(x,p) = (x',p)$  and  $\mathcal{M}_P(x,p) = (x,p')$ 

... then quantify the direct disturbance on  $x_S$ 



We define a quantifier of  $\mathcal{M}_X(x,p)$  for the second measurement

Clumsiness parameter

$$\mathbb{J}(\mathcal{M}_2) = \sum_{Q_1, Q_3 = \pm 1} \left| \Pr(Q_1, Q_3)_{\mathcal{S}_{(1, 2', 3)}} - \Pr(Q_1, Q_3)_{\mathcal{S}_{(1, 3)}} \right|$$

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# A proposed test in atomic ensembles



#### Modified LG inequality

 $\langle Q_7 Q_3 \rangle + \langle Q_5 Q_3 \rangle + \langle Q_7 Q_5 \rangle + 1 + \mathbb{I}_{37} + \mathbb{I}_{57} \ge 0$ 

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System+Meter are atoms  $\vec{J} = (N_A, J_y, J_z)$  and light  $\vec{S} = (\frac{N_L}{2}, S_y, S_z)$  $S_y^{(\text{out})} = S_y^{(\text{in})} + \kappa J_z^{(\text{in})}$ 

**noise** can be parameterized with  $\chi = \exp(-\eta N_L)$ 

$$(\Delta J_z)^2_{(\text{out})} = \chi^2 (\Delta J_z)^2_{(\text{in})} + \chi (1-\chi) \frac{N_A}{2} + (1-\chi) \frac{2}{3} N_A$$

#### Summary

- We addressed the clumsiness loophole of MR tests by
  - Exploiting the features of QND measurement
  - Oefining an appropriate clumsiness quantifier
- We extended current MR tests to exclude larger set of theories, including some invasivity
- We showed that **adroit** MR tests are feasible in **macroscopic systems** (e.g. atomic ensembles) with current-state technology

#### THANKS FOR YOUR ATTENTION!

- C. Budroni, GV, G. Colangelo, et al., PRL 115, 200403 (2015)
- GV, PhD Thesis, arxiv:1511.08104
- + in preparation

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