



Detection and chemical tailoring of entanglement in antiferromagnetic spin clusters

February 4th 2015, Bilbao

- **Molecular Nanomagnets**

 - Spin Hamiltonians

- **Entanglement detection**

 - Exchange energy as an entanglement witness

- **Spin-pair entanglement modulation**

 - Homometallic and heterometallic rings : Cr_8 , Cr_7M and Cr_{2n}M_2

 - Exchanged coupled dimers

- **Multi-spin entanglement**

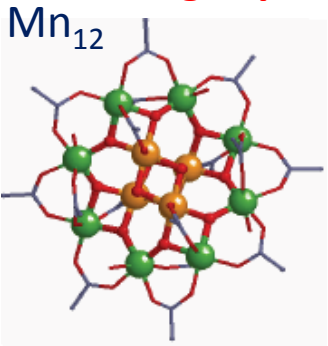
 - Energy as a witness of multi-spin entanglement

 - Homometallic and heterometallic rings

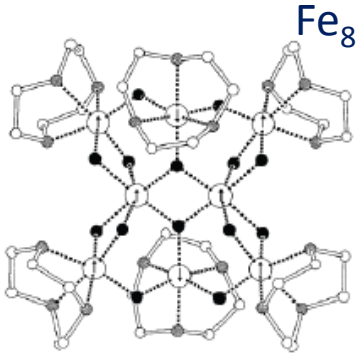
- **Conclusions**

Molecular Spin Clusters: Single Molecule Magnets

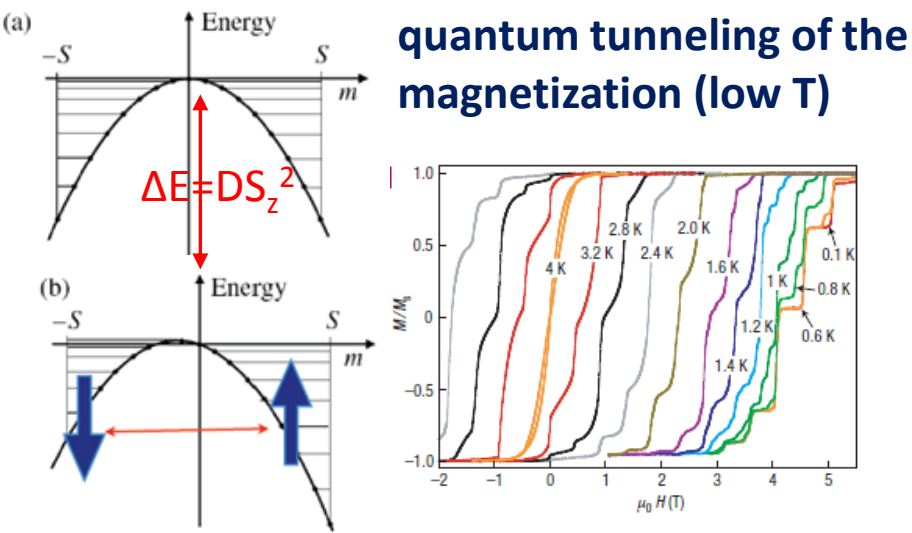
high spin ground state



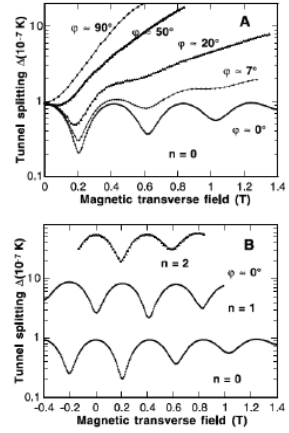
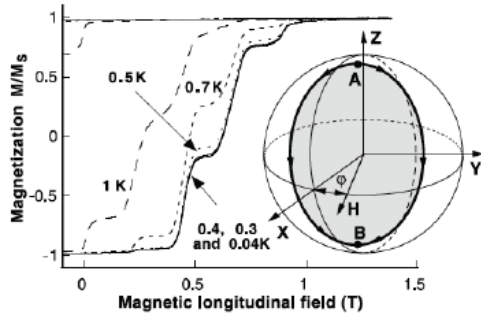
8 Mn³⁺ + 4 Mn⁴⁺



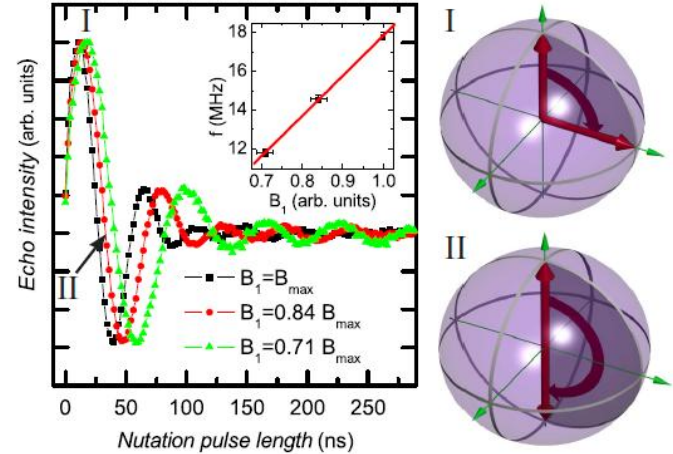
8 Mn³⁺ + 4 Mn⁴⁺



quantum phase interference

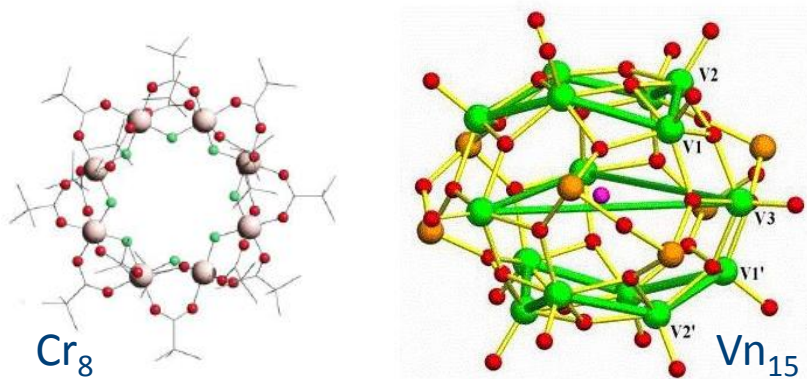


quantum oscillations



Thomas et al, Nature 383, 145 (1996).
 Sangregorio et al, Phys. Rev. Lett. 78, 4645 (1997).
 Schelgel et al., Phys. Rev. Lett. 101, 147203 (2003).

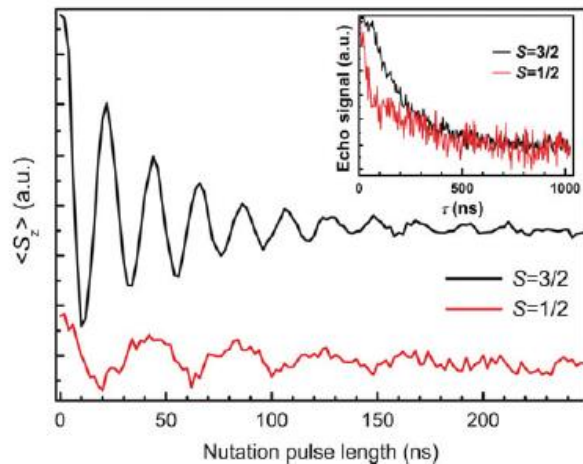
Molecular Spin Clusters: AF molecules



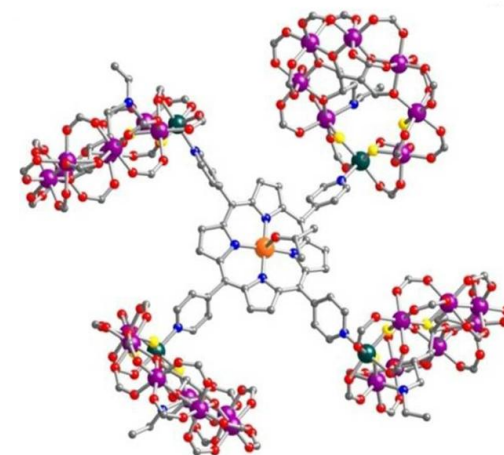
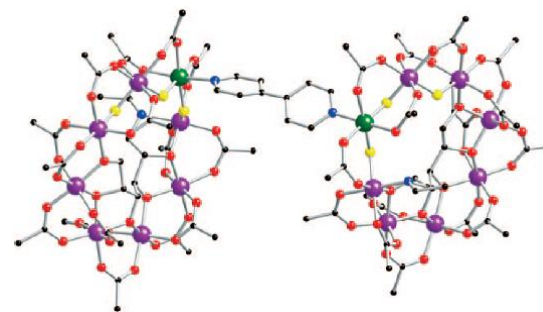
highly correlated

low spin ground state

quantum oscillations



chemical engineering



Bertaina et al, Nature 453, 203 (1996).

Timco et al. Angew. Chem. Int. Ed. 47, 9681–9684 (2008).

Timco et al., Chem. Soc. Rev. 40, 3067–3075 (2011)

Spin Hamiltonians

$$H = \sum_{i=1}^N J_{i,i+1} \mathbf{s}_i \cdot \mathbf{s}_{i+1} + \sum_{i=1}^N \left(d_i s_{i,z}^2 + \frac{1}{2} e_i (s_{i,+}^2 - s_{i,-}^2) \right) + \sum_{i < j} \mathbf{s}_i \cdot \bar{D}_{ij} \cdot \mathbf{s}_j$$

- Direct **diagonalization** of the Hamiltonian (dimension $(2s_i+1)^N$)

Irreducible Tensor Operator (ITO) formalism

block matrix (exploiting symmetries, conservation of S and M)

non-local basis: successive coupling scheme is adopted

$$| S_1 S_2 (\tilde{S}_2) S_3 (\tilde{S}_3) \dots S_{N-1} (\tilde{S}_{N-1}) S_N SM \rangle = | (\tilde{S}) SM \rangle$$

- Reduced **density matrix** at **finite temperature** in the local basis

$$\rho = \frac{1}{Z} e^{-H/k_B T} \quad | m_1, m_2, \dots, m_N \rangle$$

Entanglement Detection

An observable W is an **entanglement witness** if:

$$\langle W \rangle \geq W^{th} \text{ for all separable states } \rho_s$$

$$W^{th} = \inf_{\psi \text{ is separable}} \langle \psi | W | \psi \rangle$$

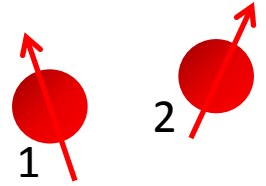
- ✓ knowledge of ρ is not required
- ✓ macroscopic properties are informative of microscopic quantum correlations
- ✓ thermal entanglement

- ✓ each witness is selective for a specific kind of entanglement

Exchange energy as a witness in AF systems

$$H = J\mathbf{s}_1 \cdot \mathbf{s}_2$$

Factorizable state



$$|\psi\rangle = |\varphi_1\rangle \otimes |\varphi_2\rangle \Rightarrow \langle H \rangle = J \langle \varphi_1 | \mathbf{s}_1 | \varphi_1 \rangle \cdot \langle \varphi_2 | \mathbf{s}_2 | \varphi_2 \rangle \geq -Js^2$$

M.R. Dowling, A. C. Doherty and S. D. Bartlett, Phys. Rev. A 70, 062113 (2004).

G. Toth, Phys. Rev. A 71, 010301 (2005).

C. Brukner and V. Vedral, arXiv quant-ph:0406040 (2004).

Exchange energy as a witness in AF systems

$$H = J\mathbf{s}_1 \cdot \mathbf{s}_2$$

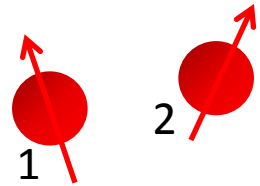
Factorizable state

$$|\psi\rangle = |\varphi_1\rangle \otimes |\varphi_2\rangle \implies \langle H \rangle \geq -Js^2$$

$$\langle H \rangle < -Js^2 \implies \text{spin-pair entanglement}$$

Singlet

$$|\psi\rangle_s \implies \langle H \rangle = -Js^2 - Js$$



M.R. Dowling, A. C. Doherty and S. D. Bartlett, Phys. Rev. A 70, 062113 (2004).

G. Toth, Phys. Rev. A 71, 010301 (2005).

C. Brukner and V. Vedral, arXiv quant-ph:0406040 (2004).

Exchange energy as a witness in AF systems

$$H = J\mathbf{s}_1 \cdot \mathbf{s}_2$$

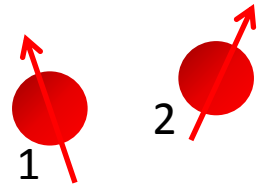
Factorizable state

$$|\psi\rangle = |\varphi_1\rangle \otimes |\varphi_2\rangle \implies \langle H \rangle \geq -Js^2$$

$$\langle H \rangle < -Js^2 \implies \text{spin-pair entanglement}$$

Singlet

$$|\psi\rangle_s \implies \langle H \rangle = -Js^2 - Js$$



M.R. Dowling, A. C. Doherty and S. D. Bartlett, Phys. Rev. A 70, 062113 (2004).

G. Toth, Phys. Rev. A 71, 010301 (2005).

C. Brukner and V. Vedral, arXiv quant-ph:0406040 (2004).

Exchange energy as a witness in AF systems

$$H = J\mathbf{s}_1 \cdot \mathbf{s}_2$$

Factorizable state

$$|\psi\rangle = |\varphi_1\rangle \otimes |\varphi_2\rangle \Rightarrow \langle H \rangle \geq -Js^2$$

$$\langle H \rangle < -Js^2 \Rightarrow \text{spin-pair entanglement}$$

Singlet

$$|\psi\rangle_s \Rightarrow \langle H \rangle = -Js^2 - Js$$

$$H = J \sum_{i=1}^N \mathbf{s}_i \cdot \mathbf{s}_{i+1}$$



$$\langle H \rangle < -JNs^2 \Rightarrow \text{spin-pair entanglement}$$

M.R. Dowling, A. C. Doherty and S. D. Bartlett, *Phys. Rev. A* 70, 062113 (2004).

G. Toth, *Phys. Rev. A* 71, 010301 (2005).

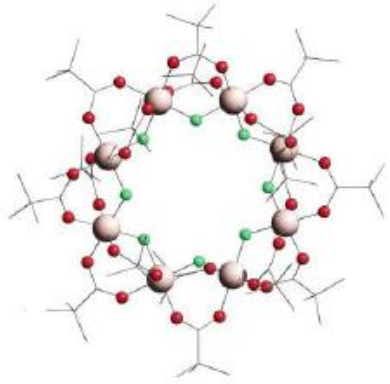
C. Brukner and V. Vedral, *arXiv quant-ph:0406040* (2004).

Spin pair entanglement modulation

I. Siloi and F. Troiani, Phys. Rev. B 86, 224404 (2012)

G. Lorusso, V. Corradini, A. Ghirri, R. Biagi, U. del Pennino, I. Siloi, F. Troiani, G. Timco, R.E.P. Winpenny and M. Affronte, Phys. Rev. B 86, 184424 (2012)

Homometallic rings: Cr₈



$$H = J \sum_{i=1}^N \mathbf{s}_i \cdot \mathbf{s}_{i+1}$$

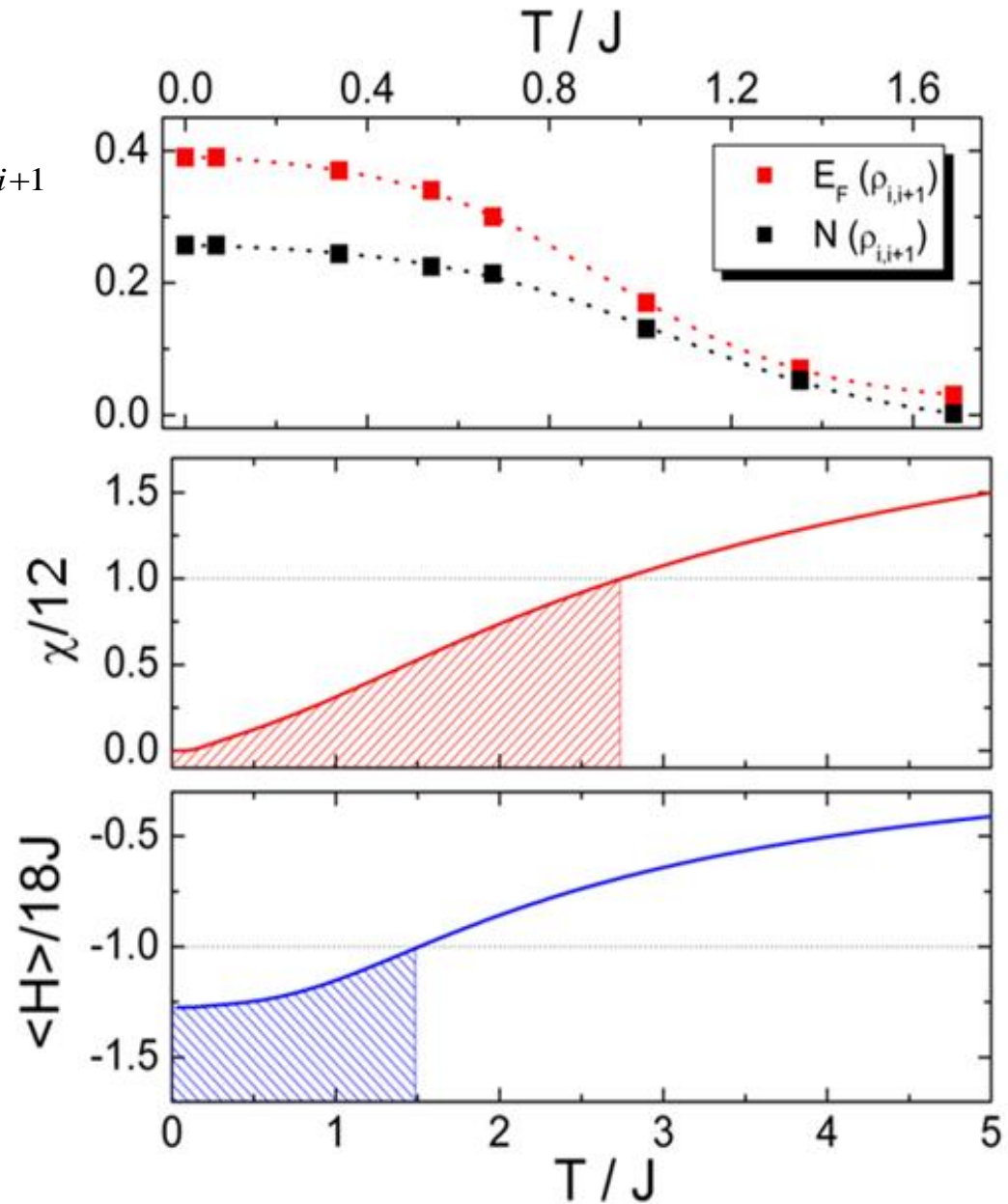
$$\Delta \approx 0.559J$$

$$J \approx 17K$$

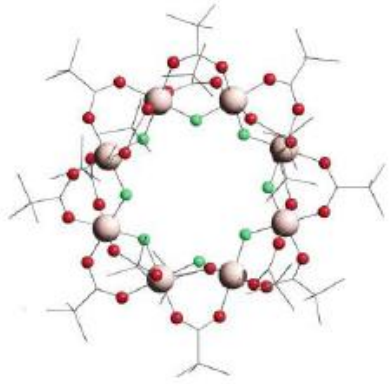
**spin-pair entanglement
@ finite temperature**

$$N(\rho_{ij}) > 0, \quad T < 1.58 J$$

$$\langle H \rangle < -18J, \quad T < 1.51 J$$



Homometallic rings: Cr₈



$$H = J \sum_{i=1}^N \mathbf{s}_i \cdot \mathbf{s}_{i+1}$$

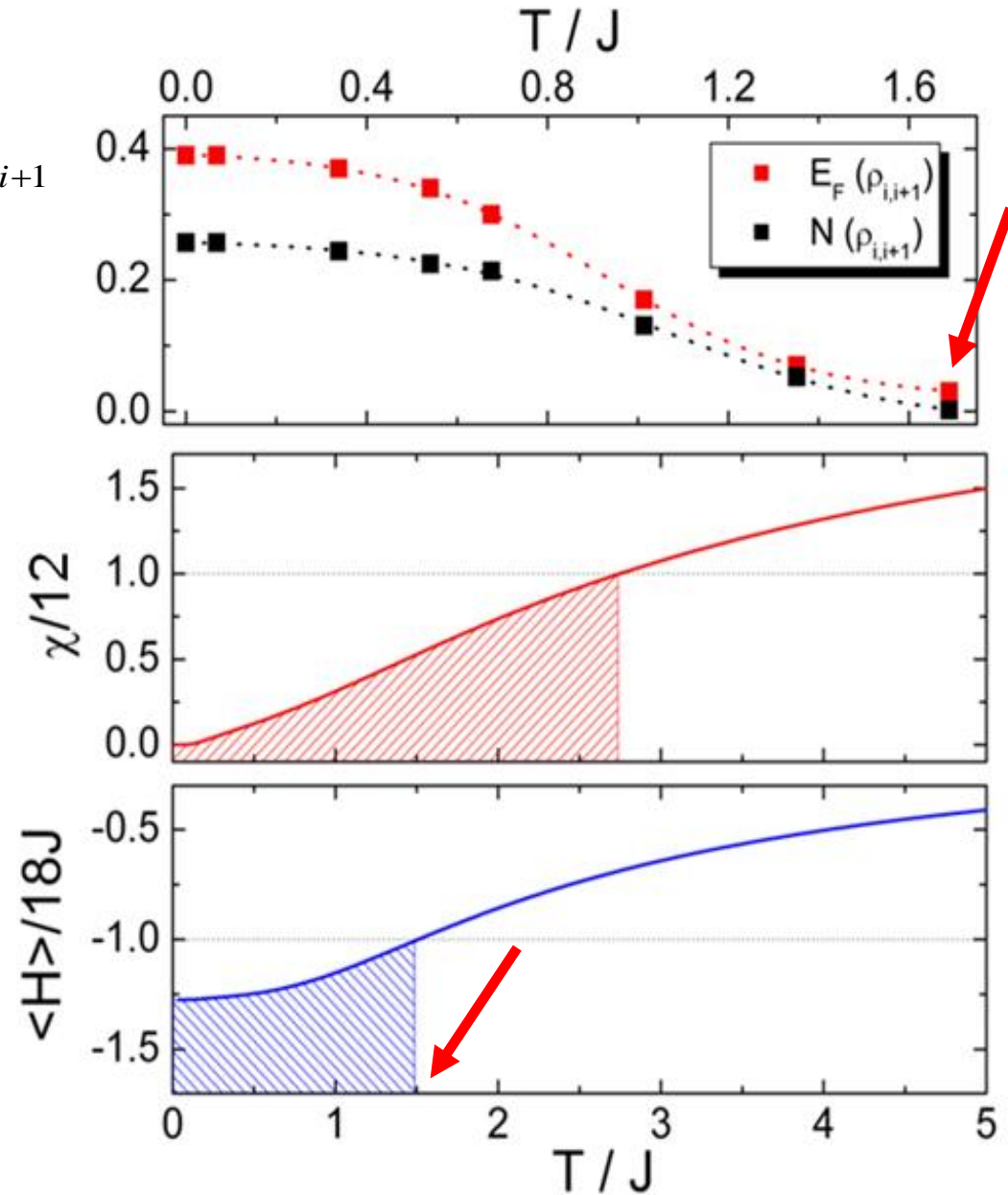
$$\Delta \approx 0.559J$$

$$J \approx 17K$$

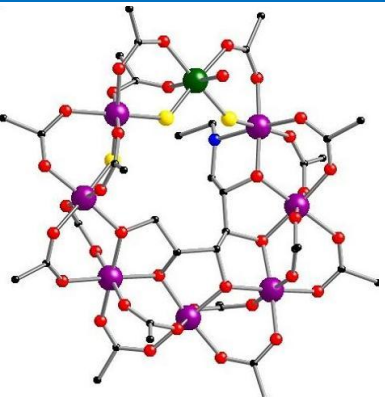
**spin-pair entanglement
@ finite temperature**

$$N(\rho_{ij}) > 0, T < 1.58 J$$

$$\langle H \rangle < -18J, T < 1.51 J$$



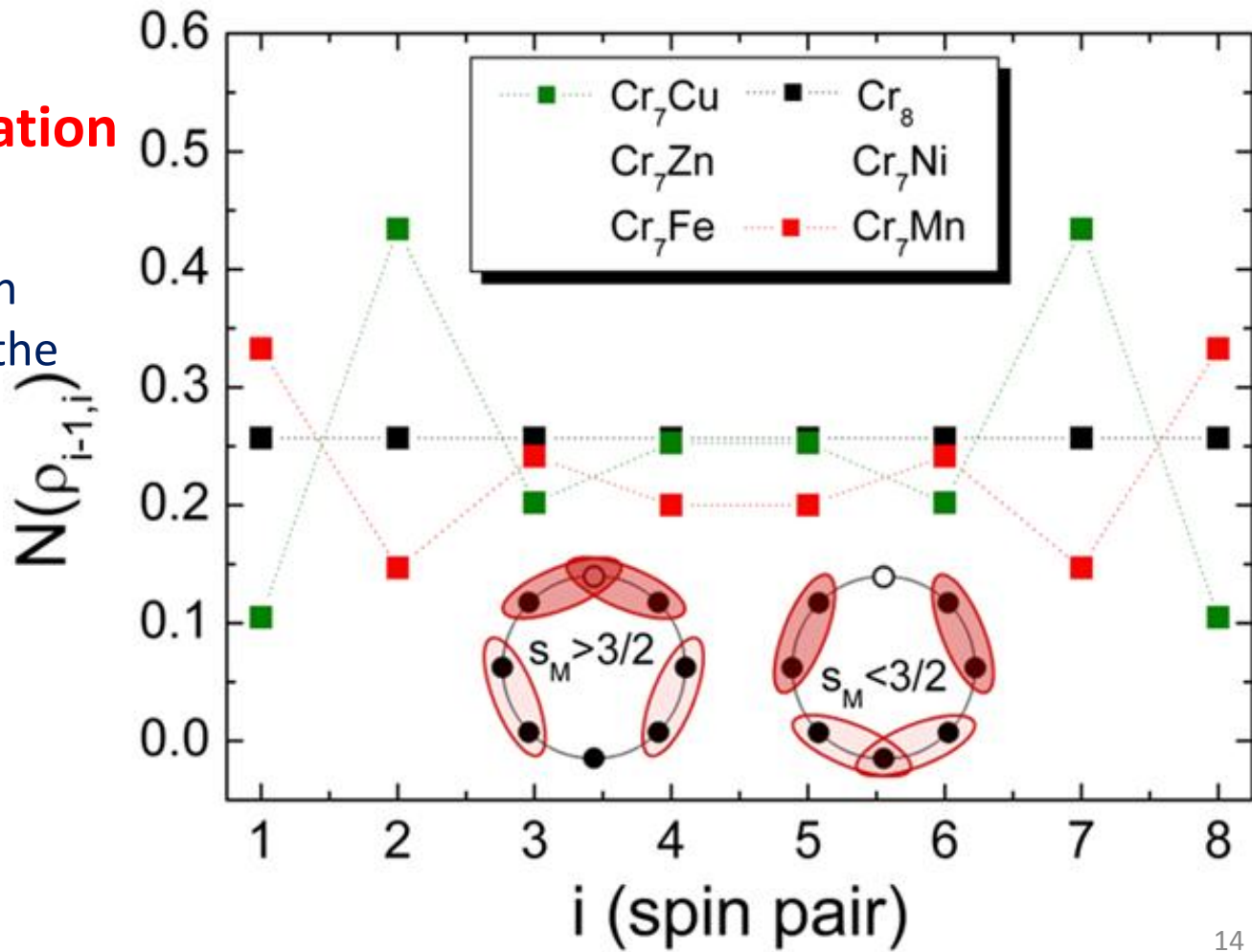
Heterometallic rings: Cr₇M



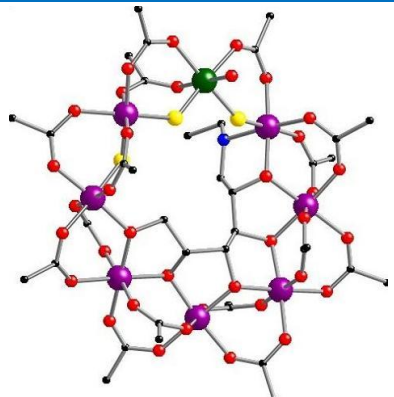
entanglement modulation @ T=0

- oscillations damped with increasing distance from the impurity
- each pair displays spin impurity dependence
- competition between neighboring pairs

	Zn	Cu	Ni	Cr	Fe	Mn
s_M	0	1/2	1	3/2	2	5/2
S	3/2	1	1/2	0	1/2	1



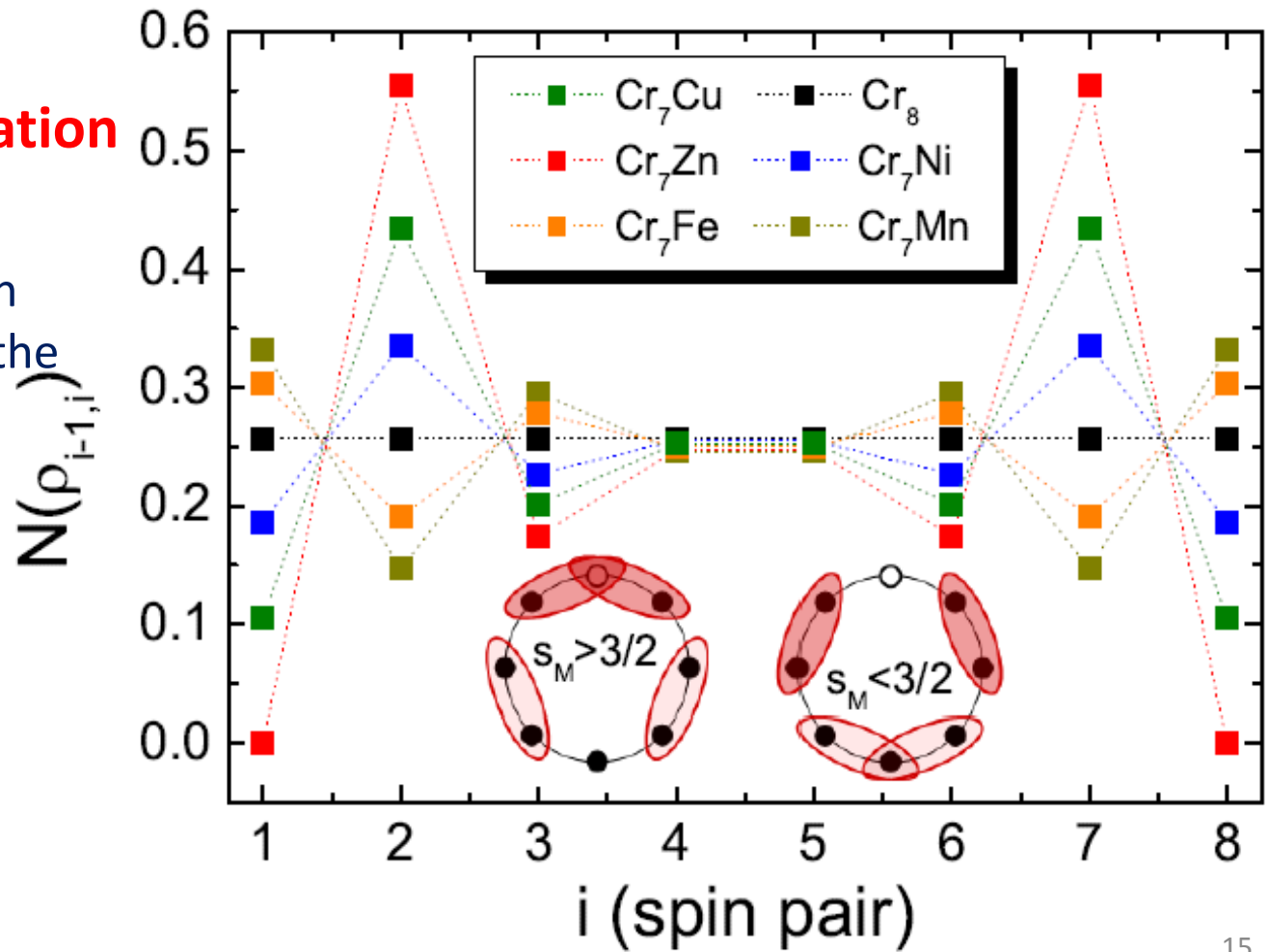
Heterometallic rings: Cr₇M



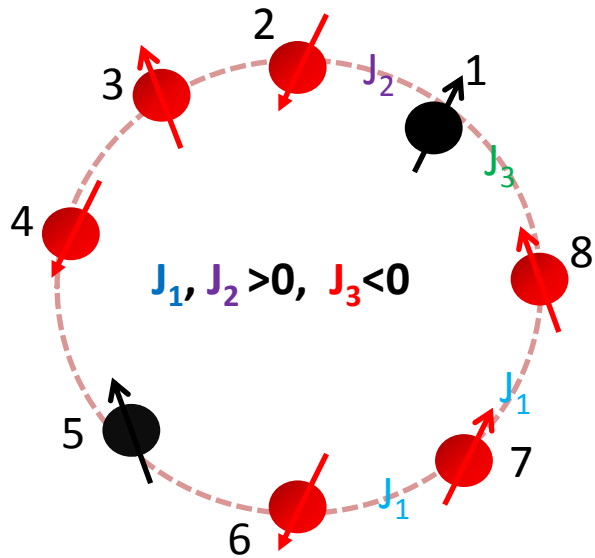
entanglement modulation @ T=0

- oscillations damped with increasing distance from the impurity
- each pair displays spin impurity dependence
- competition between neighboring pairs

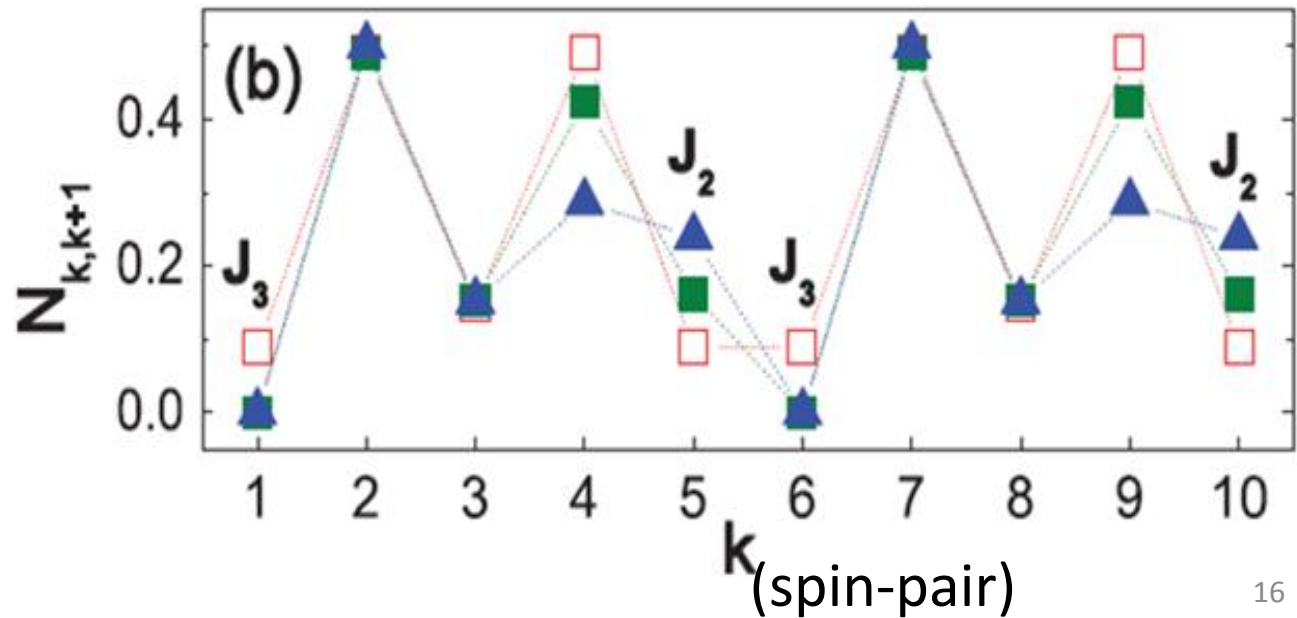
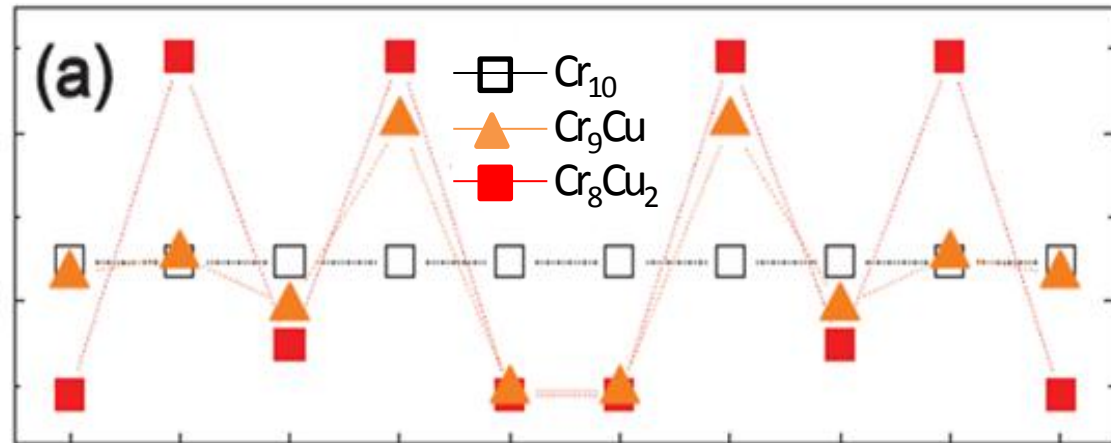
	Zn	Cu	Ni	Cr	Fe	Mn
s_M	0	1/2	1	3/2	2	5/2
S	3/2	1	1/2	0	1/2	1



Heterometallic rings: $\text{Cr}_{2n}\text{Cu}_2$



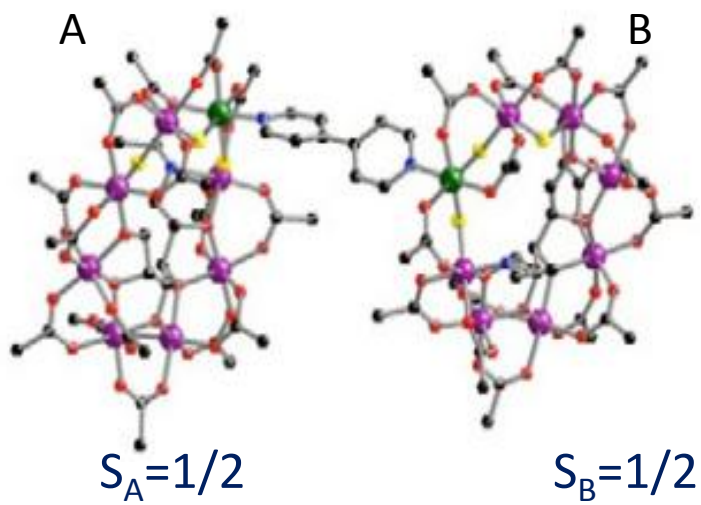
combined modulations@ T=0



- spin impurities combine modulations on spin-pair entanglement

- small antiferromagnetic coupling acts as small spin impurity

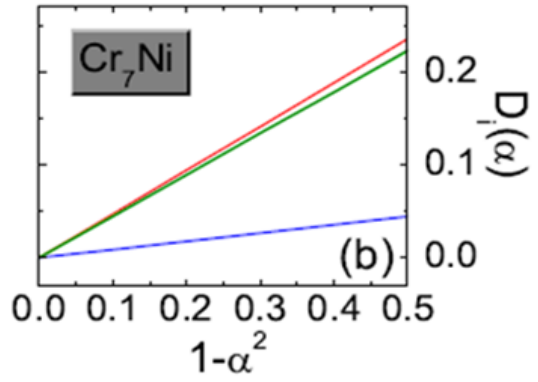
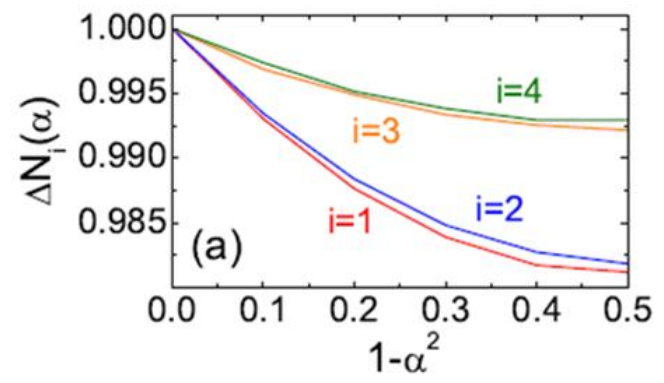
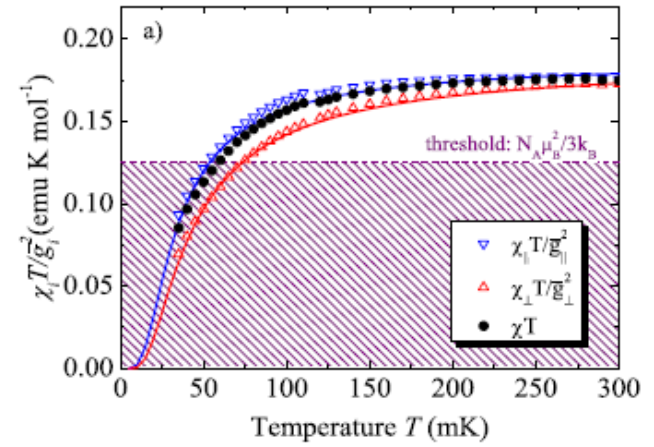
Exchanged coupled dimer



$$|\Psi_{AB}^{Ni}(\alpha)\rangle = \alpha |\uparrow\uparrow, \downarrow\downarrow\rangle - (1-\alpha^2)^{1/2} |\downarrow\downarrow, \uparrow\uparrow\rangle$$

$$\rho_{ij}^A(\alpha) = \alpha^2 \rho_{ij}^{\uparrow\uparrow} + (1-\alpha^2) \rho_{ij}^{\downarrow\downarrow}$$

$$N(\sum_i p_i \rho_i) \leq \sum_i p_i N(\rho_i)$$



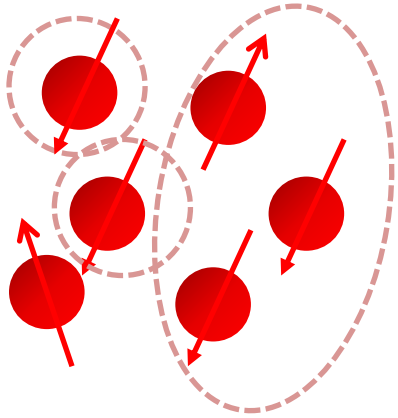
Multi-spin entanglement

F. Troiani and I. Siloi, Phys. Rev. A 86, 032330 (2012)

I. Siloi and F. Troiani, Eur. J. Phys. B (Special Issue JEMS2012) 86 (2), 1-6 (2013)

I. Siloi and F. Troiani, Phys. Rev. A 90, 042328 (2014)

k-spin entanglement

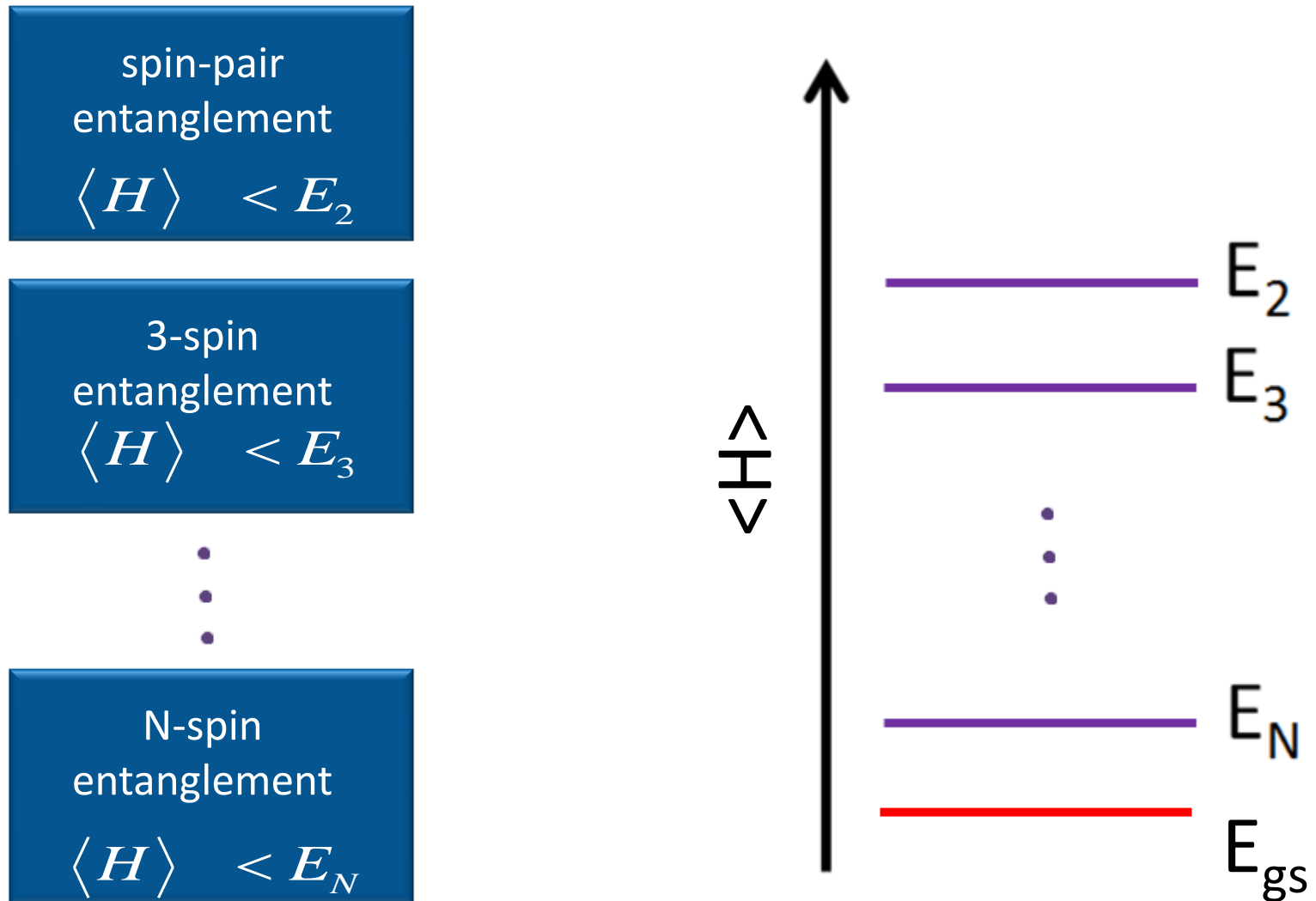


$$|\psi^k\rangle = |\varphi_{n_1}\rangle \otimes |\varphi_{n_2}\rangle \otimes \dots \otimes |\varphi_{n_p}\rangle$$

$$\sum_i n_i = N$$

$$\max\{n_i\} = k$$

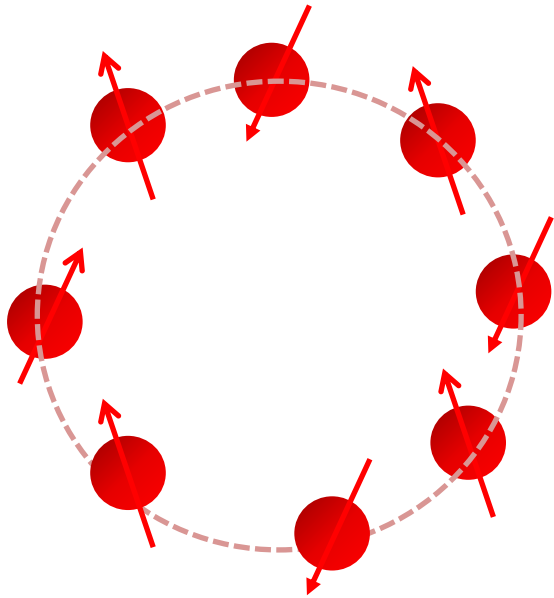
Energy as a witness of k-spin entanglement



O. Guhne, G. Toth and H. J. Briegel, New J. Phys. 7,229 (2005)

O. Guhne and G. Toth, Phys. Rev A 73, 052319 (2006)

Energy as a witness of k-spin entanglement



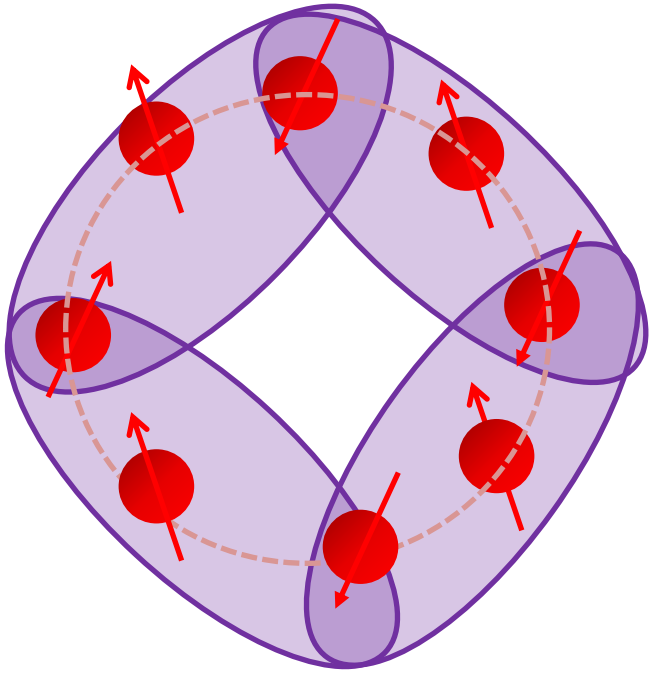
$$H = H_{123} + H_{345} + H_{567} + H_{678}$$

$$\langle H_{ijk} \rangle \geq E_3$$

fully separable or biseparable

$$\langle H \rangle > 4E_3$$

Energy as a witness of k-spin entanglement



$$H = H_{123} + H_{345} + H_{567} + H_{678}$$

$$\langle H_{ijk} \rangle \geq E_3$$

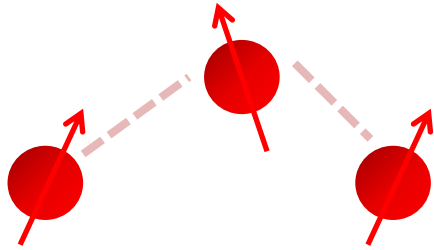
fully separable or biseparable

$$\langle H \rangle > 4E_3$$

3-spin entanglement

$$\langle H \rangle < 4E_3$$

Example: 3-spin entanglement



$$H = \mathbf{s}_1 \cdot \mathbf{s}_2 + \mathbf{s}_2 \cdot \mathbf{s}_3$$

$$E^{bs} = \langle \Phi_{bs} | H | \Phi_{bs} \rangle = \langle \varphi_1 | \mathbf{s}_1 | \varphi_1 \rangle \cdot \langle \varphi_{23} | \mathbf{s}_2 | \varphi_{23} \rangle + \langle \varphi_{23} | \mathbf{s}_2 \cdot \mathbf{s}_3 | \varphi_{23} \rangle$$

biseparable

$$|\Phi_{bs}\rangle = |\varphi_1\rangle \otimes |\varphi_{23}\rangle$$

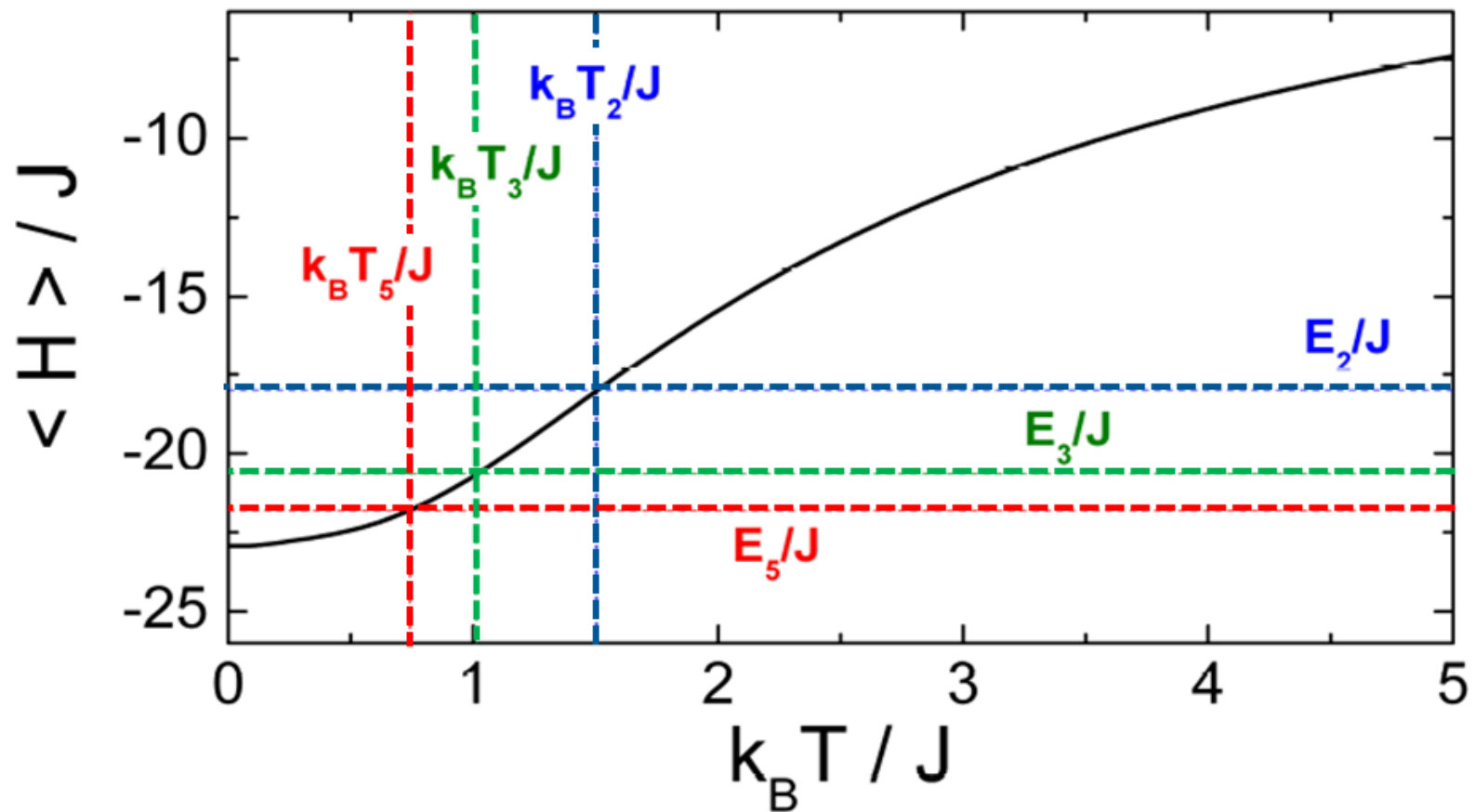
$$E^{bs} = \langle H_{eff} \rangle = \langle -s_1 s_2^z + \mathbf{s}_2 \cdot \mathbf{s}_3 \rangle$$

fully separable

$$E^{fs} = -s_2 (s_1 + s_3)$$

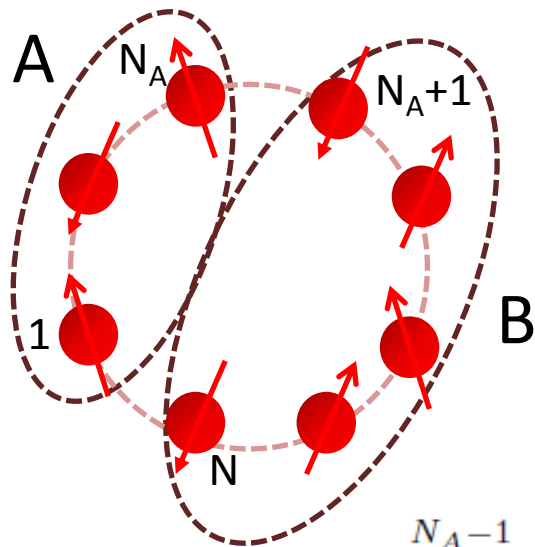
$$E < E^3 = E^{bs} = \langle -s_1 s_2^z + \mathbf{s}_2 \cdot \mathbf{s}_3 \rangle$$

Homometallic ring: Cr₈



- spin-pair entanglement $T < T_2 = 1.51$ J
- 3-spin entanglement $T < T_3 = 1.02$ J
- 5-spin entanglement $T < T_5 = 0.72$ J

k-spin entanglement in spin rings



$$H = J \sum_{i=1}^N \mathbf{s}_i \cdot \mathbf{s}_{i+1}$$

$$|\Psi_{bs}\rangle = |\Psi_A\rangle \otimes |\Psi_B\rangle$$

$$E_{bs} = \min_{N_A, N_B} E_{bs}(N_A, N_B)$$

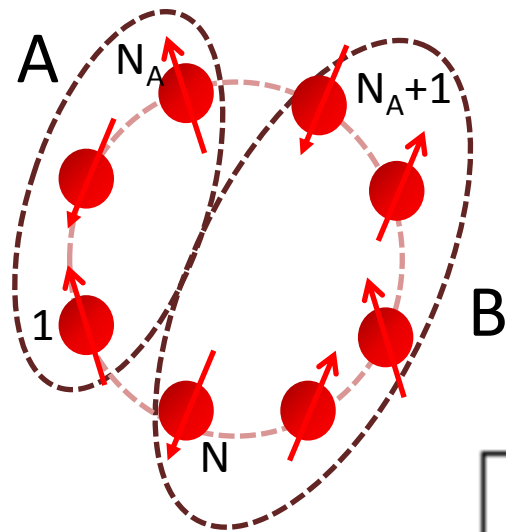
$$H = H_A + H_B + H_{AB}$$

$$H_A = \sum_{i=1}^{N_A-1} \mathbf{s}_i \cdot \mathbf{s}_{i+1}, \quad H_B = \sum_{i=N_A+1}^{N-1} \mathbf{s}_i \cdot \mathbf{s}_{i+1}, \quad H_{AB} = \mathbf{s}_{N_A} \cdot \mathbf{s}_{N_A+1} + \mathbf{s}_N \cdot \mathbf{s}_1$$

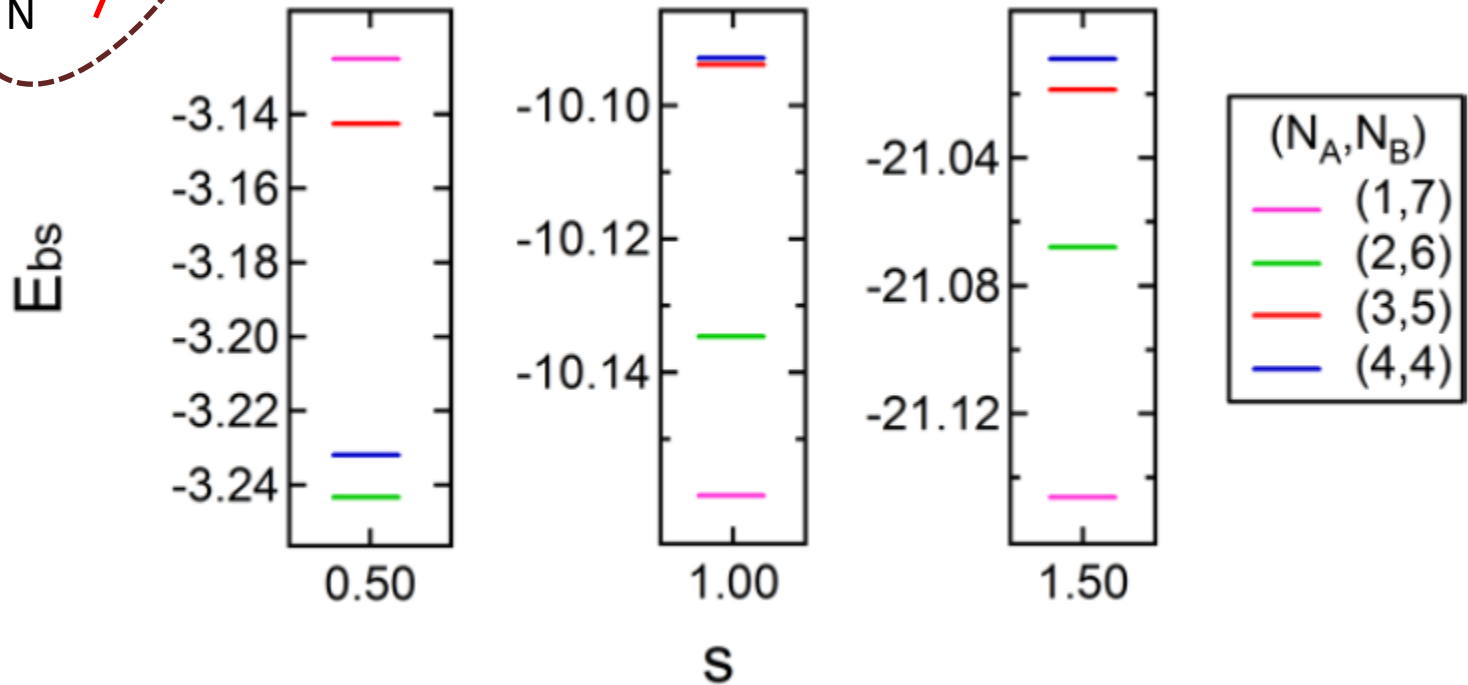
$$\begin{aligned} \langle \Psi_{bs} | H | \Psi_{bs} \rangle &= \langle \Psi_A | H_A | \Psi_A \rangle + \langle \Psi_B | H_B | \Psi_B \rangle \\ &+ \langle \Psi_A | \mathbf{s}_{N_A} | \Psi_A \rangle \cdot \langle \Psi_B | \mathbf{s}_{N_A+1} | \Psi_B \rangle \\ &+ \langle \Psi_A | \mathbf{s}_1 | \Psi_A \rangle \cdot \langle \Psi_B | \mathbf{s}_N | \Psi_B \rangle. \end{aligned}$$

$$\begin{aligned} \tilde{H}_A(z_B, z'_B) &= H_A + z_B \cdot \mathbf{s}_{N_A} + z'_B \cdot \mathbf{s}_1 \\ \tilde{H}_B(z_A, z'_A) &= H_B + z_A \cdot \mathbf{s}_{N_A+1} + z'_A \cdot \mathbf{s}_N \\ z_A &\equiv \langle \mathbf{s}_{N_A} \rangle \quad z'_A \equiv \langle \mathbf{s}_1 \rangle \\ z_B &\equiv \langle \mathbf{s}_{N_A+1} \rangle \quad z'_B \equiv \langle \mathbf{s}_N \rangle \end{aligned}$$

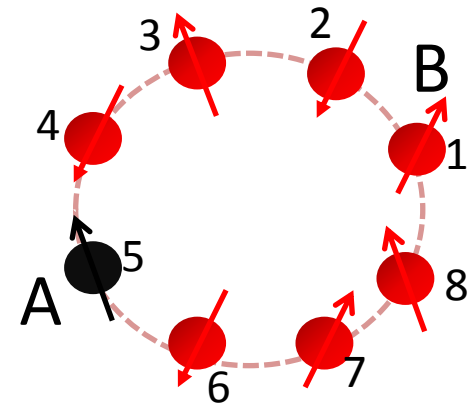
k-spin entanglement in spin rings



$$E_{bs} = \min_{N_A, N_B} E_{bs}(N_A, N_B)$$

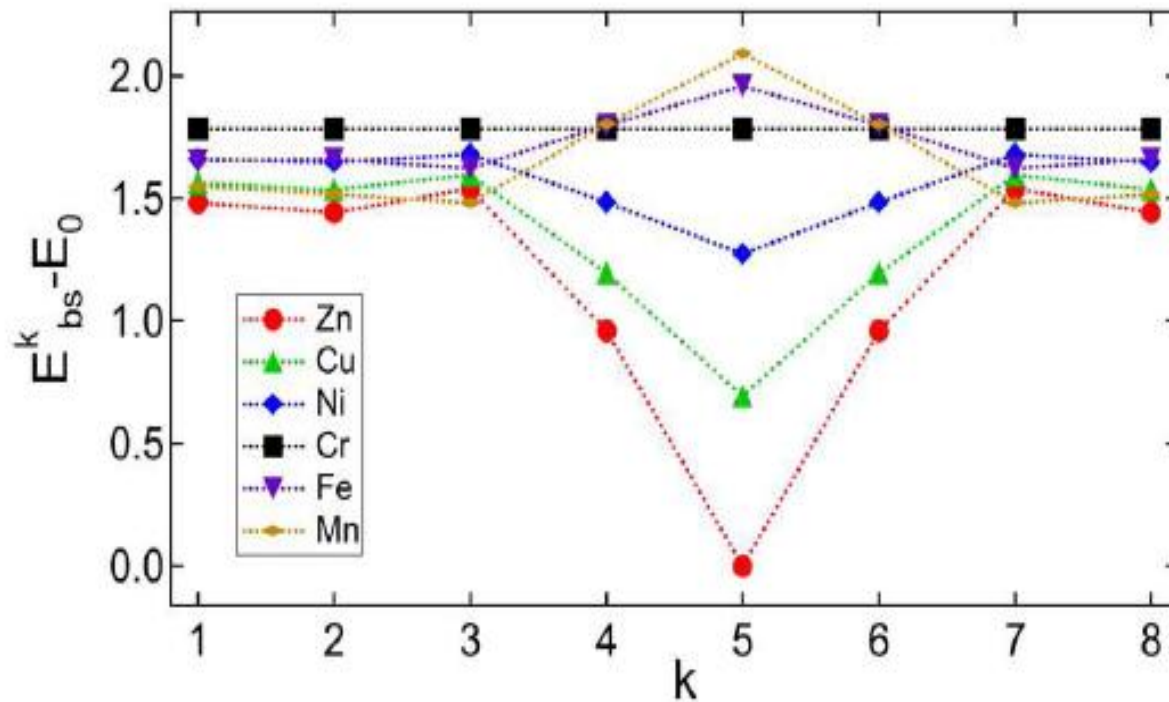


Heterometallic rings: Cr₇M

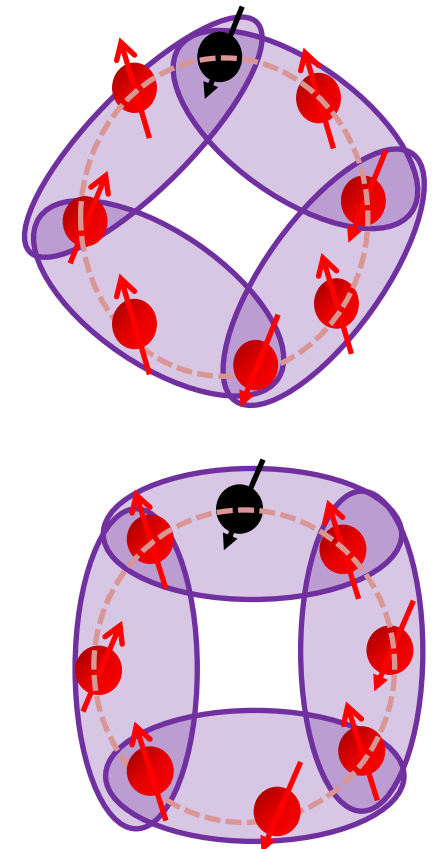


$$(N_A, N_B) = (1, 7)$$

$$\tilde{H}_B(z_A = s_k) = H_B + s_k(s_{z,k-1} + s_{z,k+1})$$



Inequivalent ways of including the defect into the 3-spin terms



Conclusions

- **chemical substitutions** induce strong **spatial modulation** of spin-pair entanglement
- **modulation** at finite temperature is detected by **spin-pair correlation function**
- **exchange** coupled **dimers**: local spins and collective spins entanglement coexist
- **exchange energy** is a witness also of **k-spin entanglement**
- **single spin separability criteria**: access to **local features** through global observable

Heterometallic rings: Cr₇M

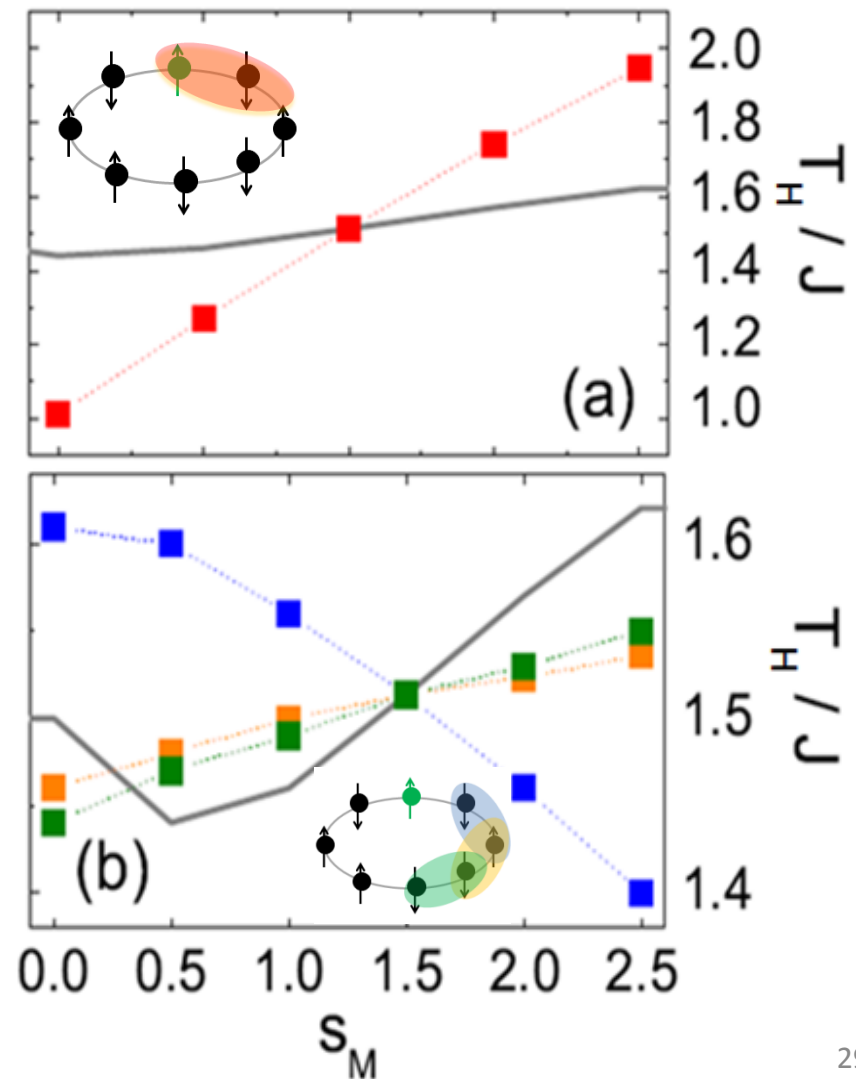
finite temperature

- spatial modulations revealed by **local witnesses**

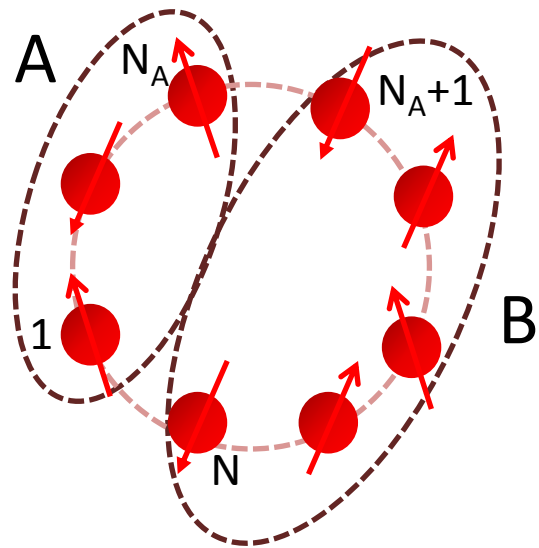
$$\langle \mathbf{S}_i \cdot \mathbf{S}_{i+1} \rangle < -S_i S_{i+1}$$

- **threshold temperatures** identified by global energy witness

$$\langle H \rangle < -(N-2)Js_{Cr}^2 - 2Js_{Cr}S_M$$



N-spin entanglement in spin rings



$$\tilde{H}_A(\mathbf{z}_B, \mathbf{z}'_B) = H_A + \mathbf{z}_B \cdot \mathbf{s}_{N_A} + \mathbf{z}'_B \cdot \mathbf{s}_1$$

$$\tilde{H}_B(\mathbf{z}_A, \mathbf{z}'_A) = H_B + \mathbf{z}_A \cdot \mathbf{s}_{N_A+1} + \mathbf{z}'_A \cdot \mathbf{s}_N$$

$$\mathbf{z}_A \equiv \langle \mathbf{s}_{N_A} \rangle \quad \mathbf{z}'_A \equiv \langle \mathbf{s}_1 \rangle$$

$$\mathbf{z}_B \equiv \langle \mathbf{s}_{N_A+1} \rangle \quad \mathbf{z}'_B \equiv \langle \mathbf{s}_N \rangle$$

$$\mathbf{z}'_A = \eta \mathbf{z}_A, \quad \mathbf{z}'_B = \eta \mathbf{z}_B \quad (\eta = \pm 1)$$

$$\cos \theta_\alpha \equiv \frac{\mathbf{z}_\alpha \cdot \mathbf{z}'_\alpha}{|\mathbf{z}_\alpha| |\mathbf{z}'_\alpha|} \neq (-1)^{N_\alpha+1} \quad Z_\alpha \equiv \left| |\mathbf{z}_\alpha| - |\mathbf{z}'_\alpha| \right| > 0,$$

$$(\mathbf{z}_B, \mathbf{z}'_B) \xrightarrow{\tilde{H}_A} (\bar{\mathbf{z}}_A, \bar{\mathbf{z}}'_A) \xrightarrow{\tilde{H}_B} (\bar{\bar{\mathbf{z}}}_B, \bar{\bar{\mathbf{z}}}'_B)$$