## k-stretchability of entanglement,

 and the duality of $k$-separability and $k$-producibility Seminar of the Department of Theoretical Physics, UPV/EHU, Bilbao
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## Introduction

Bipartite correlation and entanglement

- classification/qualification/quantification: (S)LO(CC)
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## Permutation invariant properties

- three-level structure, Young-diagrams
- $k$-partitionability ( $k$-separability), $k$-producibility (ent. depth), duality
- $k$-stretchability


## (1) Introduction

## (2) Bipartite correlation and entanglement

## (3) Multipartite correlation and entanglement

(4) Permutation symmetric properties
(5) Summary

## Quantum states

States of discrete finite quantum systems

- state vector: $|\psi\rangle \in \mathcal{H}$ (normalized) superposition
- pure state: $\pi=|\psi\rangle\langle\psi| \in \mathcal{P}$ we are uncertain about the outcomes of the measurement, pure states encode the probabilities of those


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## Mixedness and distinguishability

Measure of mixedness:

- von Neumann entropy: $S(\varrho)=-\operatorname{Tr} \varrho \ln \varrho$
- concave, nonnegative, vanishes iff $\varrho$ pure
- Schur-concavity: entropy $=$ mixedness
- increasing in bistochastic quantum channels
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Measure of distinguishability:
- (Umegaki's) quantum relative entropy: $D(\varrho \| \sigma)=\operatorname{Tr} \varrho(\ln \varrho-\ln \sigma)$
- jointly convex, nonnegative, vanishes iff $\varrho=\omega$
- quantum Stein's lemma: relative entropy $=$ distinguishability

$$
\begin{aligned}
& \text { (rate of decaying of the probability of error } \\
& \text { in hypothesis testing, Hiai \& Petz) }
\end{aligned}
$$

- decreasing in quantum channels


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Notions of correlation:

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- in q.m. there are many (nontrivially) different observables in a system
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- the state is uncorrelated iff $\operatorname{COV}(A, B)=0$ for all $A, B$, iff $\langle A B\rangle=\langle A\rangle\langle B\rangle$ for all $A, B$, iff $\varrho=\varrho_{1} \otimes \varrho_{2}$, iff $\Gamma=0$


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- state of subsystem (e.g., $\operatorname{Tr}_{2} \pi \in \mathcal{D}_{1}$ ) not necessarily pure
- $\pi$ is entangled if (and only if) $\operatorname{Tr}_{2} \pi$ and $\operatorname{Tr}_{1} \pi$ are mixed In this case, "the best possible knowledge of the whole does not involve the best possible knowledge of its parts." (Schrödinger)


## Bipartite correlation and entanglement

Mixed states: correlation

- uncorrelated: $\Gamma=0$ (product), $\varrho=\varrho_{1} \otimes \varrho_{2} \in \mathcal{D}_{\text {unc }}$, else correlated ( $\mathcal{D} \backslash \mathcal{D}_{\text {unc }}$ )
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- classically correlated sources produce states of this kind (Werner) preparable by Local Operations and Classical Communication (LOCC), else entangled ( $\mathcal{D} \backslash \mathcal{D}_{\text {sep }}$ )
- the decomposition is not unique
- deciding separability is difficult


## Bipartite correlation and entanglement - measures

- correlation "of the state itself": $\Gamma:=\varrho-\varrho_{1} \otimes \varrho_{2}$ then $\operatorname{COV}(\varrho ; A, B)=\langle A B\rangle-\langle A\rangle\langle B\rangle=\operatorname{Tr} \Gamma A \otimes B=\langle\Gamma \mid A \otimes B\rangle_{\mathrm{HS}}$
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- for the latter one, we have another, stronger motivation:

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- correlation might not be seen well from COV, but for all $A, B$,

$$
\frac{1}{2} \operatorname{CoV}(\varrho ; \hat{A}, \hat{B})^{2} \leq C(\varrho), \quad \hat{A}=A /\|A\|_{\infty}, \hat{B}=B /\|B\|_{\infty}
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- faithful: $C(\varrho)=0 \Leftrightarrow \varrho \in \mathcal{D}_{\text {unc }}, E(\varrho)=0 \Leftrightarrow \varrho \in \mathcal{D}_{\text {sep }}$
- $E(\varrho)$ is hard to calculate


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## 4 Permutation symmetric properties

## Multipartite correlation and entanglement - structure

Level 0.: subsystems
Boolean lattice structure: $P_{0}=2^{L}$

- whole system: $L=\{1,2, \ldots, n\}$
- subsystem: $X \subseteq L$, then $\mathcal{H}_{X}, \mathcal{P}_{X}, \mathcal{D}_{X}$

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Szalay, Barcza, Szilvási, Veis, Legeza, SciRep 7, 2237 (2017)
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- $\circ$


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- $\xi$-uncorrelated states: $\mathcal{D}_{\xi \text {-unc }}=\left\{\bigotimes_{X \in \xi} \varrho x\right\}$ LO-closed

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- $\xi$-separable states: $\mathcal{D}_{\xi \text {-sep }}=\operatorname{Conv} \mathcal{D}_{\xi \text {-unc }}$ LOCC-closed

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LO-monotone (proper correlation measure)

- $\xi$-entanglement (of formation):

$$
E_{\xi}(\pi)=C_{\xi} \mid \mathcal{P}(\pi), \quad E_{\xi}(\varrho)=\min \left\{\sum_{i} p_{i} E_{\xi}\left(\pi_{i}\right) \mid \sum_{i} p_{i} \pi_{i}=\varrho\right\}
$$

LOCC-monotone (proper entanglement measure)

## Multipartite correlation and entanglement - measures

Level I.: partitions lattice structure: $P_{\mathrm{I}}=\Pi(L)$

- $\xi$-correlation ( $\xi$-mutual information):

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- faithful: $C_{\xi}(\varrho)=0 \Leftrightarrow \varrho \in \mathcal{D}_{\xi \text {-unc }}, E_{\xi}(\varrho)=0 \Leftrightarrow \varrho \in \mathcal{D}_{\xi \text {-sep }}$


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## Multipartite correlation and entanglement - structure

Level II.: multiple partitions lattice structure: $P_{\mathrm{II}}=\mathcal{O}_{\downarrow}\left(P_{\mathrm{I}}\right) \backslash\{\emptyset\}$

- partition ideal: $\boldsymbol{\xi}=\left\{\xi_{1}, \xi_{2}, \ldots, \xi_{|\xi|}\right\} \subseteq P_{1}$, closed downwards w.r.t. $\preceq$
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Szalay, Barcza, Szilvási, Veis, Legeza, SciRep 7, 2237 (2017)
Szalay, PRA 92, 042329 (2015)
Szalay, Kökényesi, PRA 86, }032341\mathrm{ (2012)
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- 0

- 0


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$\xi$-separable states: $\mathcal{D}_{\xi-\text {-sep }}=\operatorname{Conv} \mathcal{D}_{\xi \text {-unc }}$

$$
\boldsymbol{v} \preceq \boldsymbol{\xi} \Leftrightarrow \mathcal{D}_{v \text {-unc }} \subseteq \mathcal{D}_{\xi-\text {-unc }}
$$ LOCC-closed

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```
Szalay, Barcza, Szilvási, Veis, Legeza, SciRep 7, 2237 (2017)
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$$

- spec.: $k$-partitionable and $k$-producible (chains)

$$
\boldsymbol{\mu}_{k}=\left\{\mu \in P_{\mathbf{I}}| | \mu \mid \geq k\right\}, \quad \boldsymbol{\nu}_{k}=\left\{\nu \in P_{\mathbf{I}}|\forall N \in \nu:|N| \leq k\}\right.
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\begin{aligned}
& \mu_{k}=\left\{\mu \in P_{1}| | \mu \mid \geq k\right\}, \quad \nu_{k}=\left\{\nu \in P_{1}|\forall N \in \nu:|N| \leq k\}\right. \\
& n=2 \text { : } \\
& \text { - } 0
\end{aligned}
$$

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\begin{array}{ll}
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- with these:
k-partitionably and k-producibly uncorrelated
k-partitionably and k-producibly separable states
Szalay, Barcza, Szilvási, Veis, Legeza, SciRep 7, 2237 (2017)
Szalay, PRA 92, 042329 (2015)
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Szalay, Barcza, Szilvási, Veis, Legeza, SciRep 7, }2237\mathrm{ (2017)
Szalay, PRA 92, 042329 (2015)
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Szalay, Barcza, Szilvási, Veis, Legeza, SciRep 7, 2237 (2017)
Szalay, PRA 92, 042329 (2015)

## Example: Electron system of molecules

- elementary subsystems: localized atomic orbitals (Pipek-Mezey)
- "atomic split": $\alpha=\left\{A_{1}, A_{2}, \ldots, A_{|\alpha|}\right\}$ (blue)
- "bond split": $\beta=\left\{B_{1}, B_{2}, \ldots, B_{|\beta|}\right\}$ (red)


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benzene $\left(\mathrm{C}_{6} \mathrm{H}_{6}\right)$ :

$$
C_{\alpha}=29.52, C_{\beta}=2.33
$$




(in units $\ln 4$ )
Szalay, Barcza, Szilvási, Veis, Legeza, SciRep 7, 2237 (2017)

## Example: Electron system of molecules

- elementary subsystems: localized atomic orbitals (Pipek-Mezey)
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- "bond split": $\beta=\left\{B_{1}, B_{2}, \ldots, B_{|\beta|}\right\}$ (red)
cyclobutadiene $\left(\mathrm{C}_{4} \mathrm{H}_{4}\right)$ :

$$
C_{\alpha}=19.48, C_{\beta}=3.17
$$


(in units $\ln 4$ )
Szalay, Barcza, Szilvási, Veis, Legeza, SciRep 7, 2237 (2017)

## Entanglement classes

Level III: Entanglement classes lattice structure: $P_{\text {III }}=\mathcal{O}_{\uparrow}\left(P_{\mathrm{II}}\right) \backslash\{\emptyset\}$

- ideal filter: $\underline{\boldsymbol{\xi}}=\left\{\boldsymbol{\xi}_{1}, \boldsymbol{\xi}_{2}, \ldots, \boldsymbol{\xi}_{|\underline{\xi}|}\right\} \subseteq P_{\mathrm{II}}$ (closed upwards w.r.t. $\preceq$ )
- partial order: $\underline{\boldsymbol{v}} \preceq \underline{\boldsymbol{\xi}}$ def.: $\underline{\boldsymbol{v}} \subseteq \underline{\boldsymbol{\xi}}$


## Entanglement classes

Level III: Entanglen

- ideal filter: $\underline{\boldsymbol{\xi}}=$
- partial order: $\underline{\imath}$


Szalay, PRA 92, 042329

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- partial separability classes: intersections of $\mathcal{D}_{\xi \text {-sep }}$

$$
\mathcal{C}_{\underline{\xi} \text {-sep }}:=\bigcap_{\xi \notin \underline{\xi}} \overline{\mathcal{D}_{\xi-\text { sep }}} \cap \bigcap_{\xi \in \underline{\xi}} \mathcal{D}_{\xi \text {-sep }}
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Szalay, PRA 92, 042329 (2015)

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- LOCC convertibility:
if $\exists \varrho \in \mathcal{C}_{\underline{\boldsymbol{v}}}, \exists \Lambda$ LOCC map s.t. $\Lambda(\varrho) \in \mathcal{C}_{\underline{\boldsymbol{\xi}}}$ then $\underline{\boldsymbol{v}} \preceq \underline{\boldsymbol{\xi}}$

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## Entanglement classes

Level III: Entanglen

- ideal filter: $\underline{\boldsymbol{\xi}}=$
- partial order: $\underline{1}$
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Szalay, PRA 92, 042329

## Correlation classes

Level III: Corr./Ent. classes lattice structure: $P_{\mathrm{III}}=\mathcal{O}_{\uparrow}\left(P_{\mathrm{II}}\right) \backslash\{\emptyset\}$

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proven constructively for $n=3$
Han, Kye, PRA 99, 032304 (2019)

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Szalay, JPhysA 51, 485302 (2018)

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## (1) Introduction

## (2) Bipartite correlation and entanglement

(3) Multipartite correlation and entanglement
(4) Permutation symmetric properties

## Permutation symmetric correlation and entanglement

Level I.: splitting type of the system of $n$ elementary subsystems

$$
n=2:
$$



## Permutation symmetric correlation and entanglement

Level I.: splitting type of the system of $n$ elementary subsystems

$$
n=3:
$$



## Permutation symmetric correlation and entanglement

Level I.: splitting type of the system of $n$ elementary subsystems
$n=4$ :


## Permutation symmetric correlation and entanglement

Level I.: splitting type of the system of $n$ elementary subsystems

- integer partition $\hat{\xi}=\left\{x_{1}, x_{2}, \ldots, x_{|\hat{\xi}|}\right\}$ of $n$ (multiset)
$n=4$ :



## Permutation symmetric correlation and entanglement

Level I.: splitting type of the system of $n$ elementary subsystems

- integer partition $\hat{\xi}=\left\{x_{1}, x_{2}, \ldots, x_{|\hat{\xi}|}\right\}$ of $n$ (multiset) (Young diag.)
- coarser/finer: $\preceq$ partial order: $\hat{v} \preceq \hat{\xi}$ if exist $v \preceq \xi$ of those types
- this is a new partial order, $\top, \perp$, not a lattice $\hat{P}_{1}$
$n=2$ :



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Structure of $k$-partitionability and $k$-producibility

- $P_{1}$ graded lattice, gradation $=$ partitionability



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- $P_{1}$ graded lattice, gradation $=$ partitionability
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Structure of $k$-partitionability and $k$-producibility

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- note: $\preceq$ is not respected by the conjugation




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## Construction

- perm. symmetric properties, not only for perm. symmetric states


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& \uparrow \mathcal{O}_{\downarrow} \backslash\{\emptyset\} \quad \uparrow \mathcal{O}_{\downarrow} \backslash\{\theta\} \\
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## k-partitionability, $k$-producibility and $k$-stretchability

height, width and rank of a Young diagram

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\begin{aligned}
h(\hat{\xi}) & :=|\hat{\xi}| \\
w(\hat{\xi}) & :=\max \hat{\xi} \\
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## (1) Introduction

## (2) Bipartite correlation and entanglement

## (3) Multipartite correlation and entanglement

4 Permutation symmetric properties
(5) Summary

## Take home message

Notions of correlations:

- pure states of classical systems are uncorrelated (product)
- correlation in pure states is of quantum origin, this is what we call entanglement


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- correlation: "how correlated $=$ how not uncorrelated"


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These principles were applied for multipartite systems painlessly.
- general case: partitions, three-level structure
- permutation invariant case: Young diagrams, conjugation
- partitionability/producibility/stretchability: height/width/rank


## Thank you for your attention!

# Szalay, arXiv:1906.10798 [quant-ph] (submitted after the talk) Szalay, JPhysA 51, 485302 (2018) <br> Szalay, Barcza, Szilvási, Veis, Legeza, SciRep 7, 2237 (2017) <br> Szalay, PRA 92, 042329 (2015) 

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of Human Capacities


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