k-stretchability of entanglement, and the duality of *k*-separability and *k*-producibility Seminar of the Department of Theoretical Physics, UPV/EHU, Bilbao

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Permutation invariant properties

- three-level structure, Young-diagrams
- k-partitionability (k-separability), k-producibility (ent. depth), duality
- k-stretchability



2 Bipartite correlation and entanglement

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Permutation symmetric properties



States of discrete finite quantum systems

- *state vector*: $|\psi\rangle \in \mathcal{H}$ (normalized) superposition
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Measure of mixedness:

- von Neumann entropy: $S(\varrho) = -\operatorname{Tr} \varrho \ln \varrho$
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Measure of distinguishability:

- (Umegaki's) quantum relative entropy: $D(\varrho || \sigma) = \operatorname{Tr} \varrho(\ln \varrho \ln \sigma)$
- \bullet jointly convex, nonnegative, vanishes iff $\varrho=\omega$
- quantum Stein's lemma: *relative entropy* = *distinguishability* (rate of decaying of the probability of error
 - in hypothesis testing, Hiai & Petz)

• decreasing in quantum channels



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- in q.m. there are many (nontrivially) different observables in a system
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- the state is *uncorrelated* iff COV(A, B) = 0 for all A, B, iff $\langle AB \rangle = \langle A \rangle \langle B \rangle$ for all A, B, iff $\rho = \rho_1 \otimes \rho_2$, iff $\Gamma = 0$

Image: Image:

Bipartite correlation and entanglement

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- $\bullet\,$ state of subsystem (e.g., $\mathsf{Tr}_2\,\pi\in\mathcal{D}_1)$ not necessarily pure
- π is entangled if (and only if) Tr₂ π and Tr₁ π are mixed In this case, "the best possible knowledge of the whole does not involve the best possible knowledge of its parts." (Schrödinger)

Mixed states: correlation

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- the decomposition is not unique
- deciding separability is difficult

- correlation "of the state itself": $\Gamma := \varrho \varrho_1 \otimes \varrho_2$ then $\text{COV}(\varrho; A, B) = \langle AB \rangle - \langle A \rangle \langle B \rangle = \text{Tr } \Gamma A \otimes B = \langle \Gamma | A \otimes B \rangle_{\text{HS}}$
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• correlation might not be seen well from COV, but for all A, B,

$$\frac{1}{2}\operatorname{COV}(\varrho; \hat{A}, \hat{B})^2 \le C(\varrho), \qquad \hat{A} = A/\|A\|_{\infty}, \hat{B} = B/\|B\|_{\infty}$$

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- faithful: $C(\varrho) = 0 \Leftrightarrow \varrho \in \mathcal{D}_{unc}, \ E(\varrho) = 0 \Leftrightarrow \varrho \in \mathcal{D}_{sep}$
- *E*(*ρ*) is hard to calculate



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Level 0.: subsystems

Boolean lattice structure: $P_0 = 2^L$

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- whole system: $L = \{1, 2, \dots, n\}$
- subsystem: $X \subseteq L$, then \mathcal{H}_X , \mathcal{P}_X , \mathcal{D}_X

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- refinement (partial order): $v \leq \xi \text{ def.: } \forall Y \in v, \exists X \in \xi : Y \subseteq X$

 Szalay, Barcza, Szilvási, Veis, Legeza, SciRep 7, 2237 (2017)

 Szalay, PRA 92, 042329 (2015)
 Seevinck, Uffink, PRA 78, 032101 (2008)

 Szalay, Kökényesi, PRA 86, 032341 (2012)
 Dür, Cirac, Tarrach, PRL 83, 3562 (1999)

Level 0.: subsystems

Boolean lattice structure: $P_0 = 2^L$

lattice structure: $P_{\rm I} = \Pi(L)$

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- whole system: $L = \{1, 2, \dots, n\}$
- subsystem: $X \subseteq L$, then \mathcal{H}_X , \mathcal{P}_X , \mathcal{D}_X

Level I.: partitions

- partition: $\xi = \{X_1, X_2, \dots, X_{|\xi|}\} \in \Pi(L)$
- refinement (partial order): $v \leq \xi \text{ def.: } \forall Y \in v, \exists X \in \xi : Y \subseteq X$
- ξ -uncorrelated states: \mathcal{D}_{ξ -unc} = { $\bigotimes_{X \in \xi} \varrho_X$ } LO-closed $v \preceq \xi \Leftrightarrow \mathcal{D}_{v$ -unc $\subseteq \mathcal{D}_{\xi$ -unc
- ξ -separable states: \mathcal{D}_{ξ -sep} = Conv \mathcal{D}_{ξ -unc LOCC-closed $v \preceq \xi \Leftrightarrow \mathcal{D}_{v}$ -sep $\subseteq \mathcal{D}_{\xi}$ -sep

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Level I.: partitions

lattice structure: $P_{I} = \Pi(L)$

• *ξ-correlation* (*ξ*-mutual information):

$$C_{\xi}(\varrho) = \min_{\sigma \in \mathcal{D}_{\xi \text{-unc}}} D(\varrho || \sigma) = \sum_{X \in \xi} S(\varrho_X) - S(\varrho)$$

LO-monotone (proper correlation measure)

Szalay, Barcza, Szilvási, Veis, Legeza, SciRep 7, 2237 (2017) Szalay, PRA **92**, 042329 (2015)

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LO-monotone (proper correlation measure)

• ξ -entanglement (of formation):

$$E_{\xi}(\pi) = C_{\xi}|_{\mathcal{P}}(\pi), \qquad E_{\xi}(\varrho) = \min\left\{\sum_{i} p_{i} E_{\xi}(\pi_{i}) \mid \sum_{i} p_{i} \pi_{i} = \varrho\right\}$$

LOCC-monotone (proper entanglement measure)

Szalay, Barcza, Szilvási, Veis, Legeza, SciRep 7, 2237 (2017) Szalay, PRA **92**, 042329 (2015)

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LOCC-monotone (proper entanglement measure)

• faithful: $C_{\xi}(\varrho) = 0 \Leftrightarrow \varrho \in \mathcal{D}_{\xi\text{-unc}}, \ E_{\xi}(\varrho) = 0 \Leftrightarrow \varrho \in \mathcal{D}_{\xi\text{-sep}}$

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Szalay, Barcza, Szilvási, Veis, Legeza, SciRep 7, 2237 (2017) Szalay, PRA **92**, 042329 (2015)

Level II.: multiple partitions lattice structure: $P_{II} = \mathcal{O}_{\perp}(P_{I}) \setminus \{\emptyset\}$

Image: Image:

- partition ideal: $\boldsymbol{\xi} = \{\xi_1, \xi_2, \dots, \xi_{|\boldsymbol{\xi}|}\} \subseteq P_{\mathsf{I}}$, closed downwards w.r.t. \preceq
- partial order: $v \leq \xi$ def.: $v \subseteq \xi$

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• spec.: *k*-partitionable and *k*-producible (chains) $\mu_{k} = \{ \mu \in P_{I} \mid |\mu| \ge k \}, \qquad \nu_{k} = \{ \nu \in P_{I} \mid \forall N \in \nu : |N| \le k \}$ n = 2:

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Szalay, Barcza, Szilvási, Veis, Legeza, SciRep 7, 2237 (2017) Szalay, PRA **92**, 042329 (2015) Szalay, Kökényesi, PRA **86**, 032341 (2012)

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Level II.: multiple partitions • *ξ-correlation:* lattice structure: $P_{II} = \mathcal{O}_{\downarrow}(P_I) \setminus \{\emptyset\}$

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LO-monotone (proper correlation measure)

Szalay, Barcza, Szilvási, Veis, Legeza, SciRep 7, 2237 (2017) Szalay, PRA **92**, 042329 (2015)

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LOCC-monotone (proper entanglement measure)

Szalay, Barcza, Szilvási, Veis, Legeza, SciRep 7, 2237 (2017) Szalay, PRA **92**, 042329 (2015)

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Szalay, Barcza, Szilvási, Veis, Legeza, SciRep 7, 2237 (2017) Szalay, PRA **92**, 042329 (2015)

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Image: Image:

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Szalay, Barcza, Szilvási, Veis, Legeza, SciRep 7, 2237 (2017) Szalay, PRA **92**, 042329 (2015)

Example: Electron system of molecules

- elementary subsystems: localized atomic orbitals (Pipek-Mezey)
- "atomic split": $\alpha = \{A_1, A_2, \dots, A_{|\alpha|}\}$ (blue)
- "bond split": $\beta = \{B_1, B_2, ..., B_{|\beta|}\}$ (red)

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(in units ln 4)

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(in units ln 4)

Szalay, Barcza, Szilvási, Veis, Legeza, SciRep 7, 2237 (2017)

Level III: Entanglement classes lattice structure: $P_{III} = \mathcal{O}_{\uparrow}(P_{II}) \setminus \{\emptyset\}$

- ideal filter: $\underline{\xi} = \{\xi_1, \xi_2, \dots, \xi_{|\xi|}\} \subseteq P_{\mathsf{II}}$ (closed upwards w.r.t. \preceq)
- partial order: $\underline{v} \preceq \underline{\xi} \text{ def.: } \underline{v} \subseteq \underline{\xi}$

Szalay, PRA 92, 042329 (2015)

Entanglement classes



Szilárd Szalay (MTA Wigner RCP)

k-stretchability

June 20, 2019 18 / 28

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$$\mathcal{C}_{\underline{\xi}\text{-sep}} := \bigcap_{\xi \notin \underline{\xi}} \overline{\mathcal{D}_{\xi\text{-sep}}} \cap \bigcap_{\xi \in \underline{\xi}} \mathcal{D}_{\xi\text{-sep}}$$

Szalay, PRA 92, 042329 (2015)

Entanglement classes



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LOCC convertibility:
 if ∃ρ ∈ C_v, ∃Λ LOCC map s.t. Λ(ρ) ∈ C_ξ then v ≤ ξ

Szalay, PRA 92, 042329 (2015)

Entanglement classes



Correlation classes

Level III: Corr./Ent. classes

lattice structure:
$$P_{III} = \mathcal{O}_{\uparrow}(P_{II}) \setminus \{\emptyset\}$$

• partial separability classes: intersections of $\mathcal{D}_{\xi\text{-sep}}$

$$\mathcal{C}_{\underline{\xi}\operatorname{-sep}} := igcap_{\xi \notin \underline{\xi}} \overline{\mathcal{D}_{\xi\operatorname{-sep}}} \cap igcap_{\xi \in \underline{\xi}} \mathcal{D}_{\xi\operatorname{-sep}}$$

• LOCC convertibility: if $\exists \varrho \in C_{\underline{v}\text{-sep}}, \exists \Lambda \text{ LOCC map s.t. } \Lambda(\varrho) \in C_{\boldsymbol{\xi}\text{-sep}}$ then $\underline{v} \preceq \underline{\boldsymbol{\xi}}$

Szalay, PRA 92, 042329 (2015)

Correlation classes

Level III: Corr./Ent. classes lattice structure: $P_{III} = O_{\uparrow}(P_{II}) \setminus \{\emptyset\}$ • partial correlation classes: intersections of $D_{\xi-unc}$

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• partial separability classes: intersections of $\mathcal{D}_{\xi\text{-sep}}$

$$\mathcal{C}_{\underline{\xi} ext{-sep}} := igcap_{\xi
otin \underline{\xi}} \overline{\mathcal{D}_{\xi ext{-sep}}} \cap igcap_{\xi
otin \underline{\xi}
otin$$

• LOCC convertibility: if $\exists \varrho \in C_{\underline{\upsilon}\text{-sep}}, \exists \Lambda \text{ LOCC map s.t. } \Lambda(\varrho) \in C_{\underline{\xi}\text{-sep}}$ then $\underline{\upsilon} \preceq \underline{\xi}$

Szalay, PRA 92, 042329 (2015)
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Szalay, PRA 92, 042329 (2015)

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proven constructively for n = 3 Han, Kye, PRA 99, 032304 (2019)

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Szalay, PRA 92, 042329 (2015)

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Szalay, JPhysA 51, 485302 (2018)

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Szalay, PRA 92, 042329 (2015)



2 Bipartite correlation and entanglement

3 Multipartite correlation and entanglement





Level I.: splitting type of the system of *n* elementary subsystems

n = 2:



Szalay, arXiv:1906.10798 [quant-ph]

Szilárd Szalay (MTA Wigner RCP)

Level I.: splitting type of the system of *n* elementary subsystems

n = 3:



Szalay, arXiv:1906.10798 [quant-ph]

Szilárd Szalay (MTA Wigner RCP)

Level I.: splitting type of the system of *n* elementary subsystems

n = 4: Szalay, arXiv:1906.10798 [quant-ph]

Szilárd Szalay (MTA Wigner RCP)

Level I.: splitting type of the system of *n* elementary subsystems • integer partition $\hat{\xi} = \{x_1, x_2, \dots, x_{|\hat{\xi}|}\}$ of *n* (multiset)



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Structure of *k*-partitionability and *k*-producibility

• P_{I} graded lattice, gradation = partitionability



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• note: \leq is not respected by the conjugation

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k-stretchability

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height, width and rank of a Young diagram

$$h(\hat{\xi}) := |\hat{\xi}|$$
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height, width and rank of a Young diagram \implies properties

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chains

 $\hat{\mu}_{l} \preceq \hat{\mu}_{k} \quad \Longleftrightarrow \quad l \ge k \qquad \qquad \hat{\nu}_{l} \preceq \hat{\nu}_{k}, \quad \hat{\tau}_{l} \preceq \hat{\tau}_{k} \quad \Longleftrightarrow \quad l \le k$

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$$\mu \preceq \widehat{\nu}_{k}, \quad \widehat{\tau}_{l} \preceq \widehat{\tau}_{k} \iff l \le k$$

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• bounds among properties: $\hat{\mu}_k \preceq \hat{\nu}_{n+1-k}$, $\hat{\nu}_k \preceq \hat{\mu}_{\lceil n/k \rceil}$, from $\lceil n/w \rceil \leq h \leq n-w+1$ $\lceil n/h \rceil \leq w \leq n-h+1$

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$$I \preceq \hat{\boldsymbol{\nu}}_{k}, \quad \hat{\boldsymbol{\tau}}_{I} \preceq \hat{\boldsymbol{\tau}}_{k} \iff I \le k$$

$$+1-k, \quad \hat{\boldsymbol{\nu}}_{k} \preceq \hat{\boldsymbol{\mu}}_{\lceil n/k \rceil}, \text{ from}$$

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$$h(\hat{\xi}^{\dagger}) = w(\hat{\xi}), \quad w(\hat{\xi}^{\dagger}) = h(\hat{\xi}), \quad r(\hat{\xi}^{\dagger}) = -r(\hat{\xi}),$$

height, width and rank of a Young diagram \implies properties





2 Bipartite correlation and entanglement

3 Multipartite correlation and entanglement

4 Permutation symmetric properties



Take home message

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- pure states of classical systems are uncorrelated (product)
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- general case: partitions, three-level structure
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- partitionability/producibility/stretchability: height/width/rank

Szalay, arXiv:1906.10798 [quant-ph] (submitted after the talk) Szalay, JPhysA **51**, 485302 (2018) Szalay, Barcza, Szilvási, Veis, Legeza, SciRep **7**, 2237 (2017) Szalay, PRA **92**, 042329 (2015)

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