

# Entanglement detection from randomized measurements

Satoya Imai @ University of Siegen

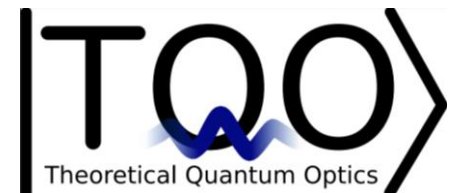
Quantum Glue Meetings in Bilbao: 16. 11. 2022

[PRL \(2021\)<sup>⊗2</sup> & PRA \(2022\)](#)

[arXiv:2205.08447 & ongoing projects](#)



DAAD



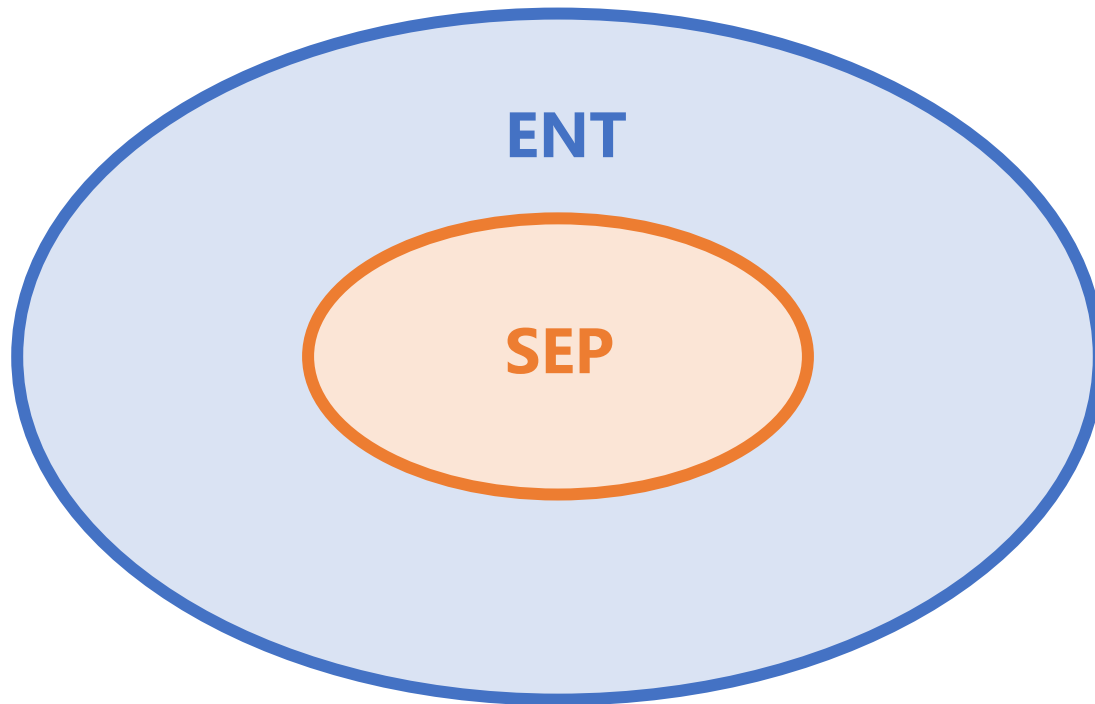




# Outline

- **Basics of entanglement**
- **Randomized measurements**
- **Conclusion**

# Basics of entanglement



# Recap: Single-qudit state

## ■ Pure state

- $d$ -dimensional complex vector  $|\psi\rangle \in H_d$ , s.t.  $\langle\psi|\psi\rangle = 1$

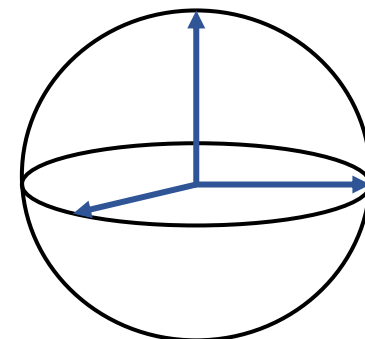
## ■ Mixed stat

- $d \times d$ -matrix  $\rho \in H_d$ , s.t.  $\rho^\dagger = \rho$  and  $\text{tr}[\rho] = 1$  and  $\rho \geq 0$

## ■ Single-qubit state: ( $d=2$ )

$$\rho = \frac{1}{2} \left( I + \sum_{i=x,y,z} \langle\sigma_i\rangle_\rho \sigma_i \right), \quad \sum \langle\sigma_i\rangle_\rho^2 \leq 1$$

Bloch ball ( $d=2$ )



# Bipartite entanglement

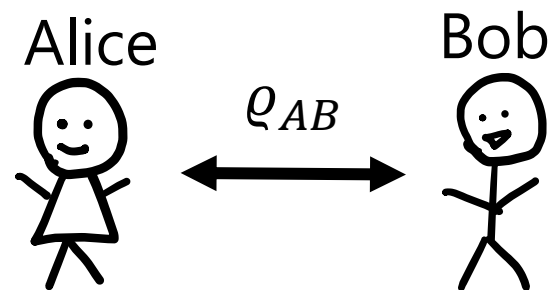
## ■ Pure state

- Product:  $|\psi\rangle_{\text{prod}} = |\phi\rangle_A \otimes |\phi\rangle_B$

- ▶  $|0\rangle_A \otimes |0\rangle_B$

- Otherwise, entangled

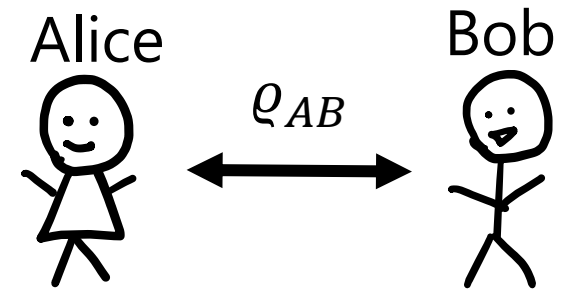
- ▶  $|\text{Bell}\rangle = \frac{1}{\sqrt{2}}(|00\rangle_{AB} + |11\rangle_{AB})$



# Bipartite entanglement

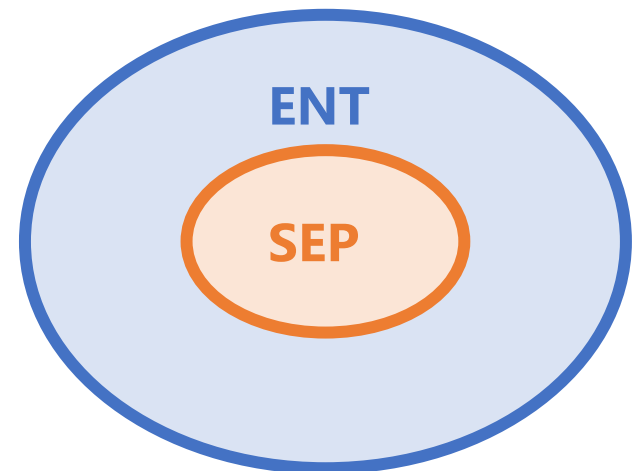
## ■ Pure state

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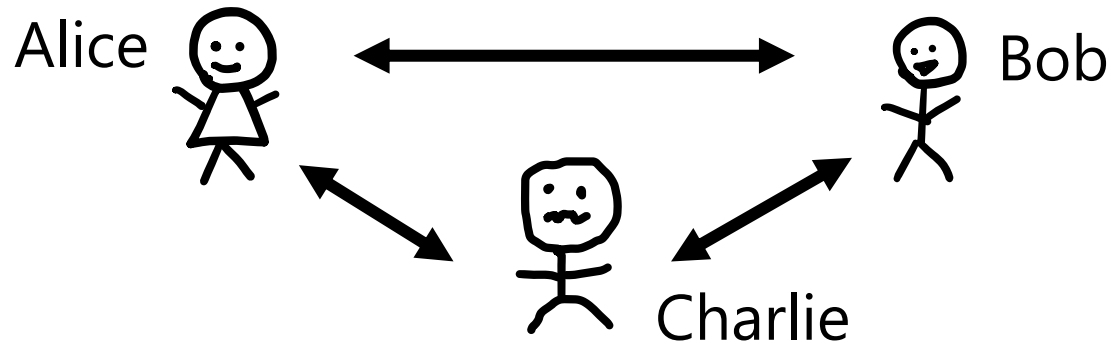


## ■ Mixed state

- Separable:  $\rho_{\text{sep}} = \sum p_i \rho_i^A \otimes \rho_i^B$
- Otherwise, entangled



# Multipartite entanglement

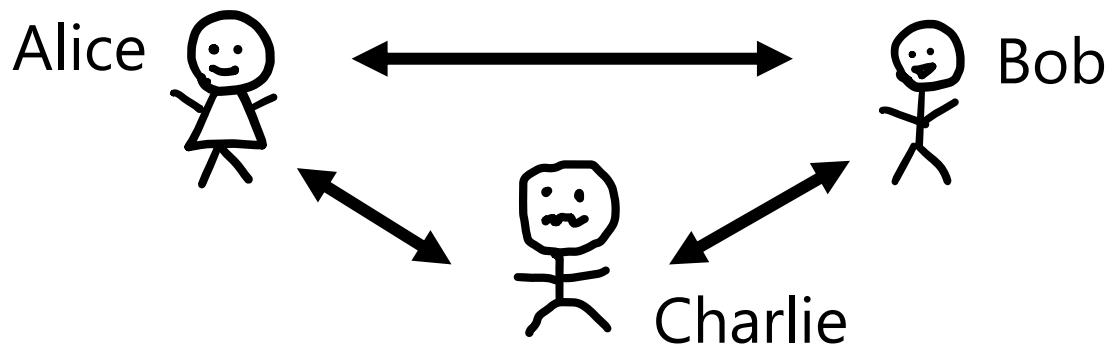


## ■ There are different forms

- Fully separable:  $|\psi\rangle_{fs} = |\phi\rangle_A \otimes |\phi\rangle_B \otimes |\phi\rangle_C$ 
  - ▶  $|0\rangle \otimes |0\rangle \otimes |0\rangle$



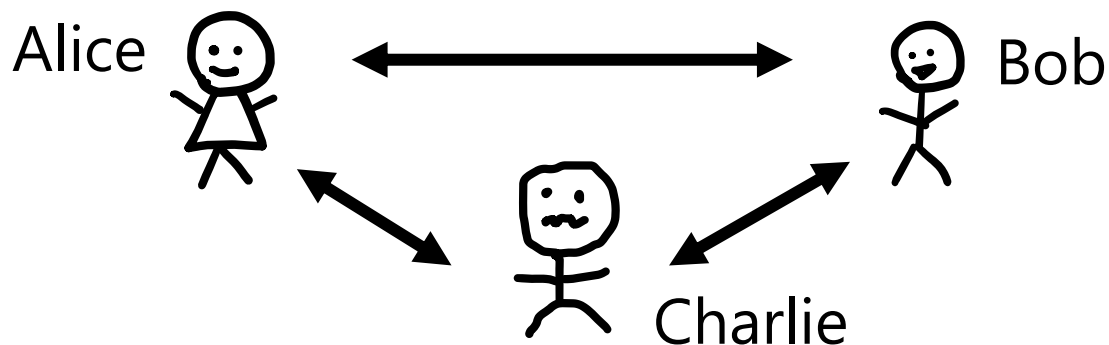
# Multipartite entanglement



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  - ▶  $|0\rangle \otimes |0\rangle \otimes |0\rangle$
- Biseparable:  $|\psi\rangle_{bs} = |\phi\rangle_{AB} \otimes |\phi\rangle_C$ 
  - ▶  $|\text{Bell}\rangle \otimes |0\rangle$

# Multipartite entanglement



## ■ There are different forms

- Fully separable:  $|\psi\rangle_{fs} = |\phi\rangle_A \otimes |\phi\rangle_B \otimes |\phi\rangle_C$ 
  - ▶  $|0\rangle \otimes |0\rangle \otimes |0\rangle$
- Biseparable:  $|\psi\rangle_{bs} = |\phi\rangle_{AB} \otimes |\phi\rangle_C$ 
  - ▶  $|\text{Bell}\rangle \otimes |0\rangle$
- Otherwise, genuine multiparticle entangled (GME)
  - ▶  $|\text{GHZ}\rangle = \frac{1}{\sqrt{2}}(|000\rangle + |111\rangle)$  or  $|\text{W}\rangle = \frac{1}{\sqrt{3}}(|001\rangle + |010\rangle + |100\rangle)$

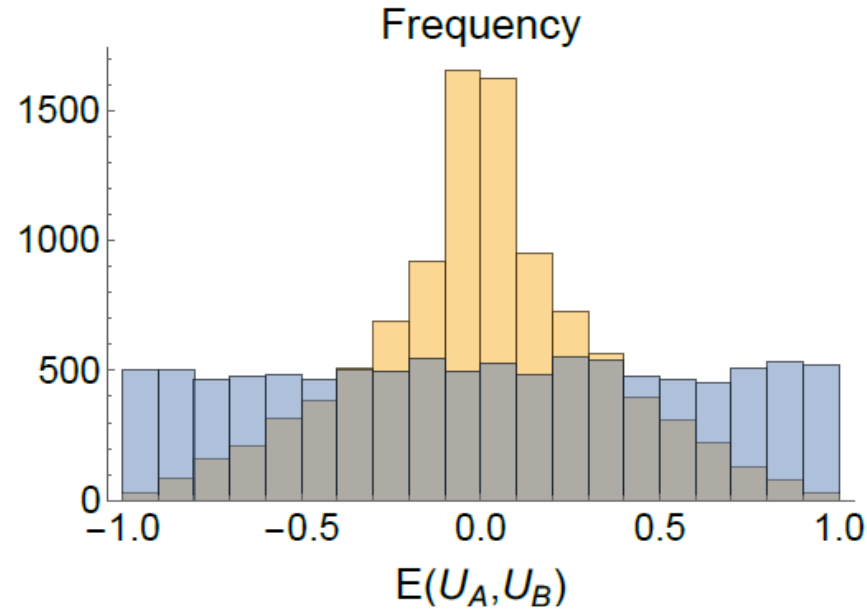
# Separability problem

**How to decide whether a given state is separable or not?**

## ■ Importance

- Experimental generation of entanglement
- Fundamental difference from classical correlations
- Mathematical decomposition of matrices

# Randomized Measurements





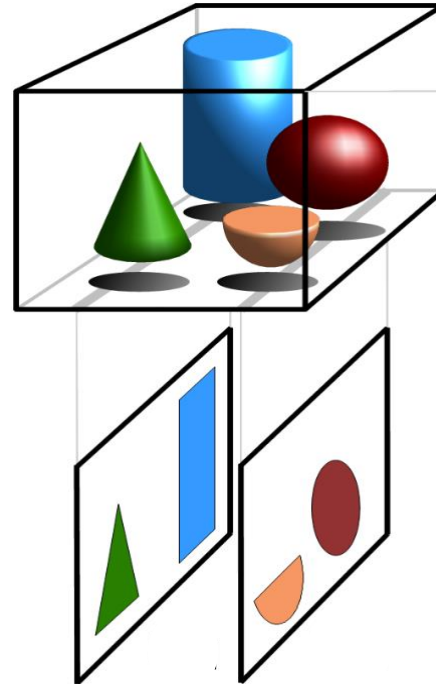
**Our approach is different from shadow tomography\***

\*HY Huang, R Kueng, J Preskill, Nature Physics 2020

# Why randomized measurements?

## ■ Several challenges

- State tomography

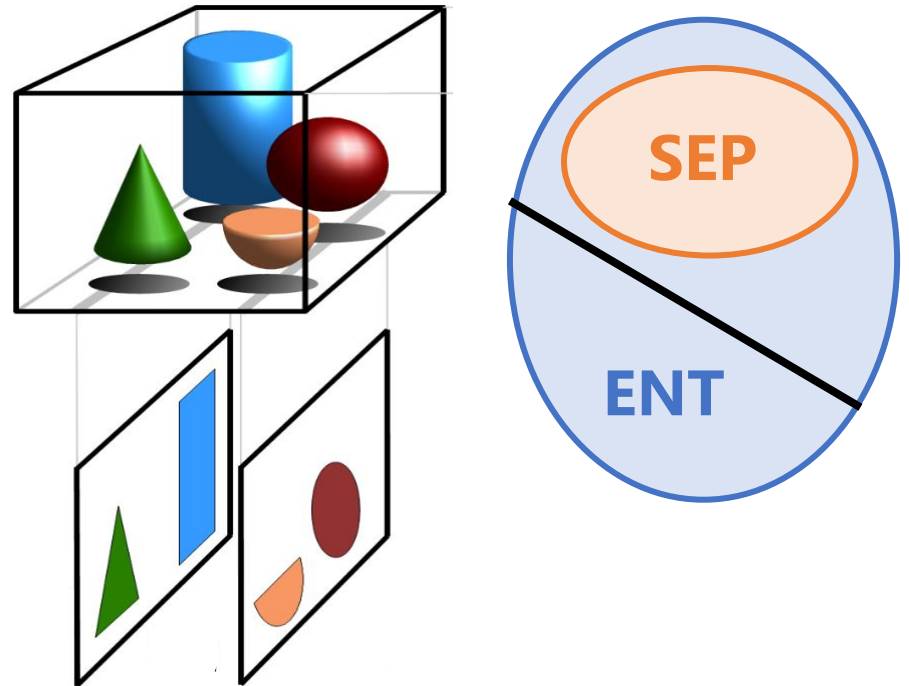




# Why randomized measurements?

## ■ Several challenges

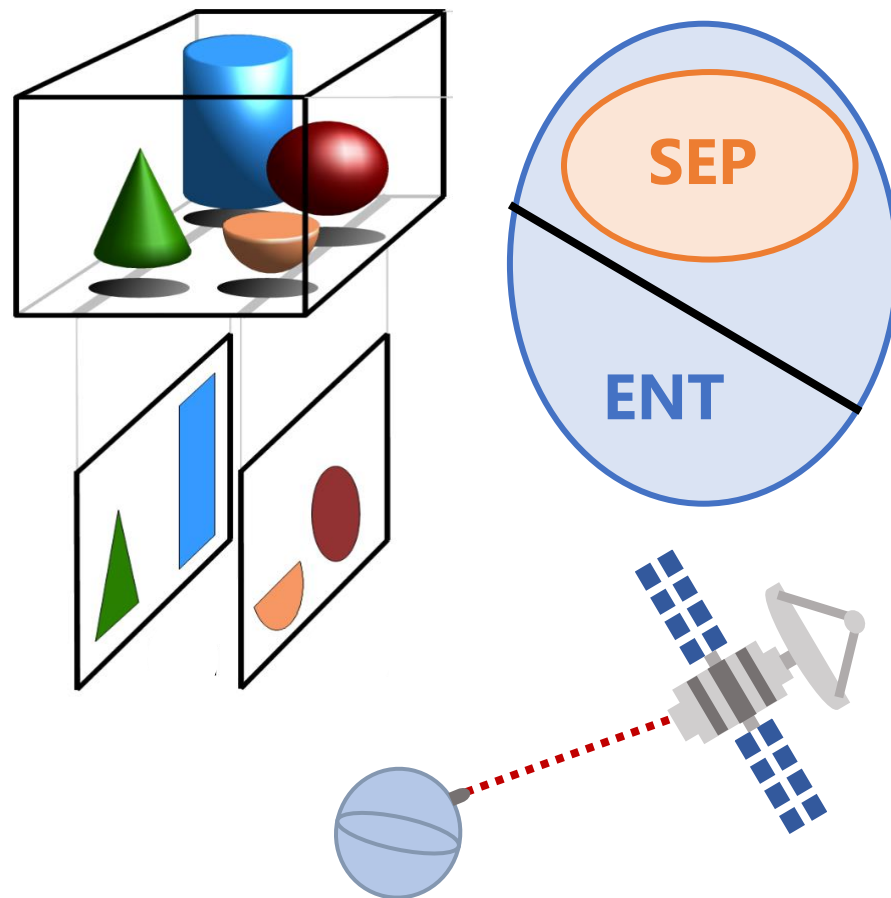
- State tomography
- Entanglement witness
- 



# Why randomized measurements?

## ■ Several challenges

- State tomography
- Entanglement witness
- Sharing reference frame

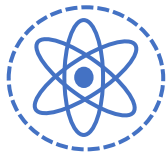


**Randomized measurements!**

# What are randomized measurements?

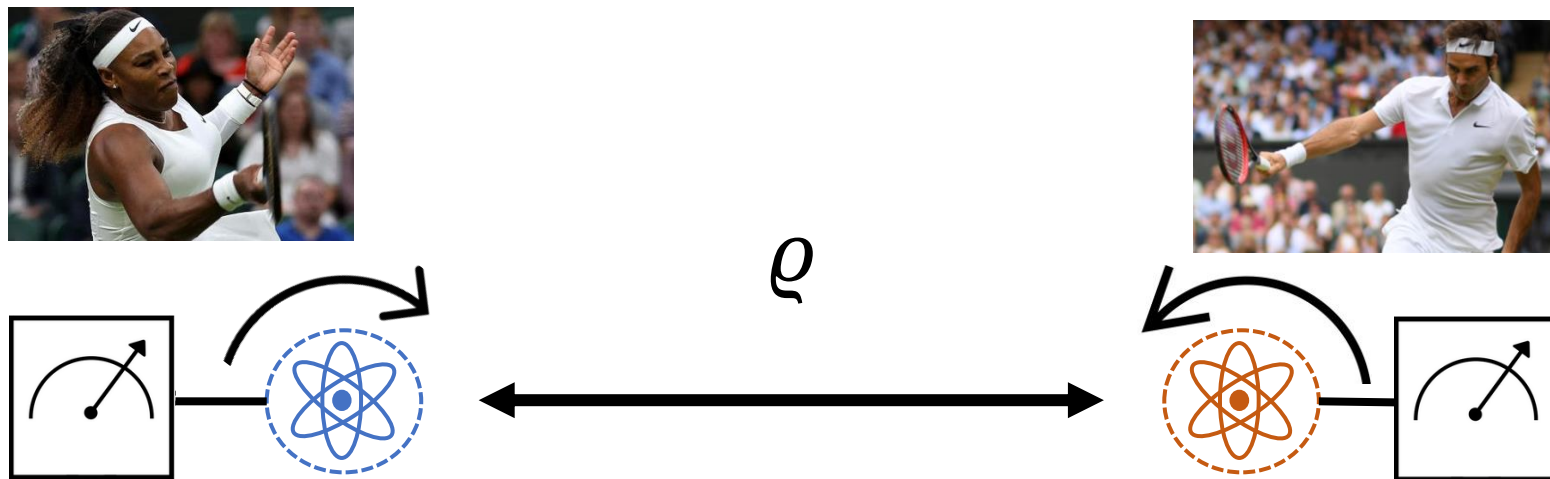


$\rho$



- They do not share common reference frame

# What are randomized measurements?

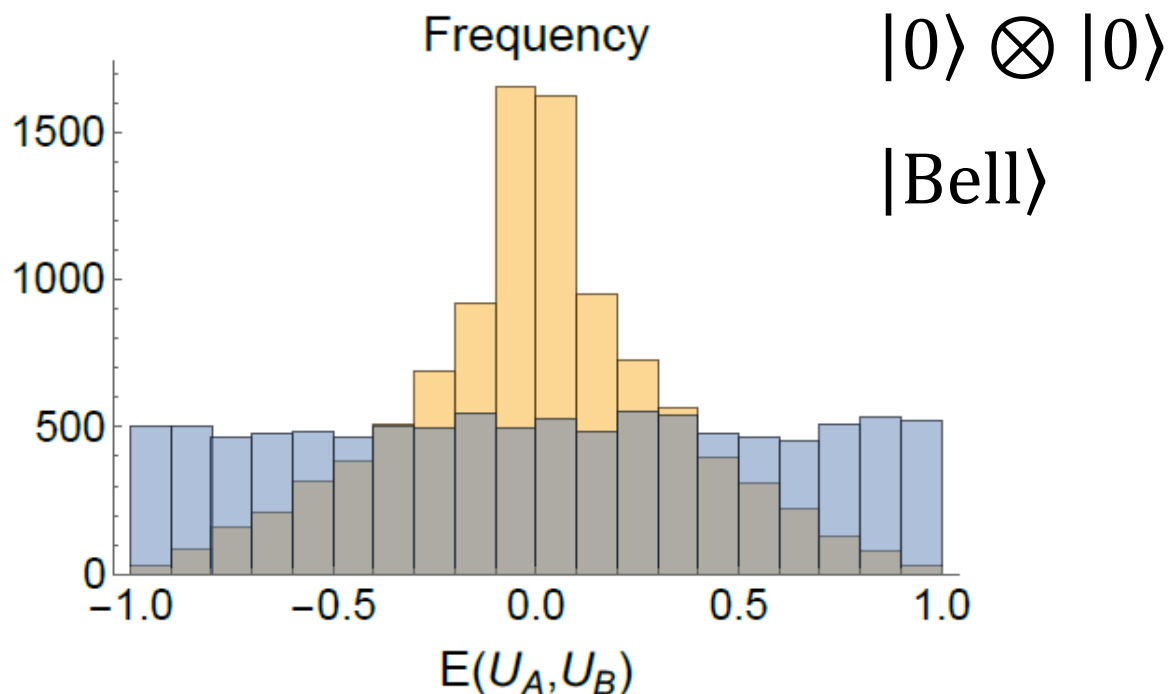


- They do not share common reference frame
- **Idea:** *Rotate* measurement direction arbitrarily

# Random correlation function

$$E(U_A, U_B) = \text{tr}(\rho_{AB} U_A^\dagger \sigma_Z U_A \otimes U_B^\dagger \sigma_Z U_B)$$

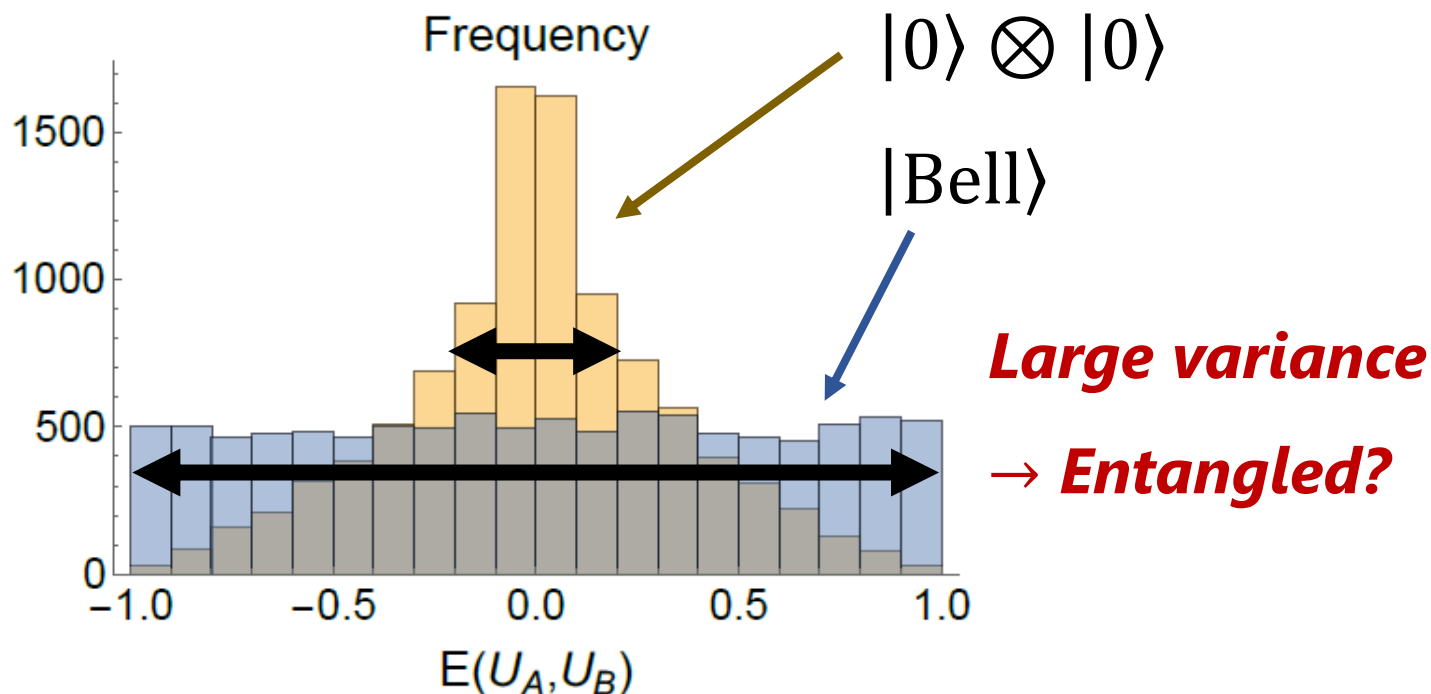
## Sample over random unitaries



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## Sample over random unitaries





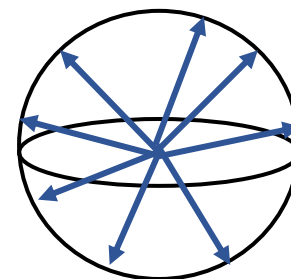
# Moments over Haar unitaries

$$\mathcal{R}_{AB}^{(r)} = N \int dU_A \int dU_B \left[ \text{tr}(\rho_{AB} U_A^\dagger \sigma_z U_A \otimes U_B^\dagger \sigma_z U_B) \right]^r$$

**Simplifications: integral  $\rightarrow$  sum**

$$\mathcal{R}_{AB}^{(2)} \stackrel{!}{=} \sum_{i,j=x,y,z} \langle \sigma_i \otimes \sigma_j \rangle_{\rho_{AB}}^2$$

Rotated Bloch vector



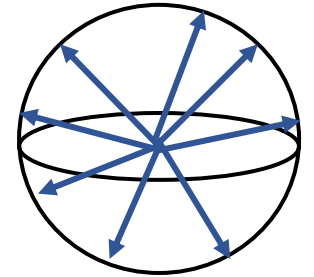
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**Simplifications: integral  $\rightarrow$  sum**

Rotated Bloch vector

$$\mathcal{R}_{AB}^{(2)} \stackrel{!}{=} \sum_{i,j=x,y,z} \langle \sigma_i \otimes \sigma_j \rangle_{\rho_{AB}}^2$$



**Reference-frame-independent entanglement detection**

$$\mathcal{R}_{AB}^{(2)} > 1 \Rightarrow \rho_{AB} \text{ is entangled}$$

# Geometry of two-qubit states

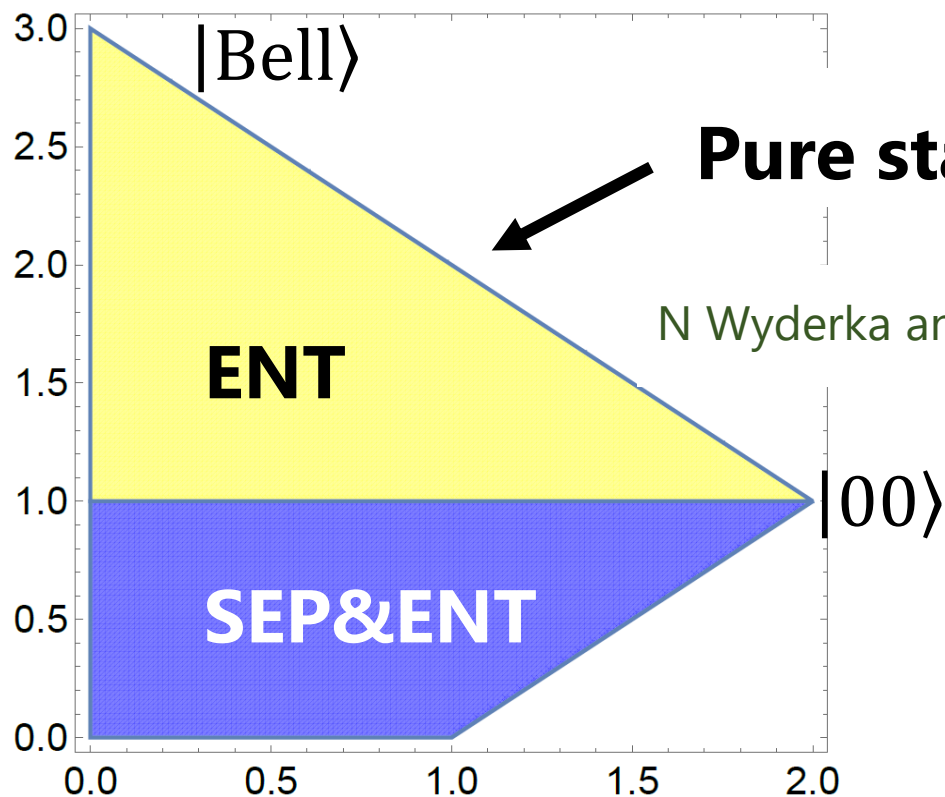
$$\mathcal{R}_{AB}^{(2)} = \sum \langle \sigma_i \otimes \sigma_j \rangle_{\rho_{AB}}^2$$



$$\mathcal{R}_A^{(2)} + \mathcal{R}_B^{(2)} = \sum \langle \sigma_i \rangle_{\rho_A}^2 + \sum \langle \sigma_i \rangle_{\rho_B}^2$$

# Geometry of two-qubit states

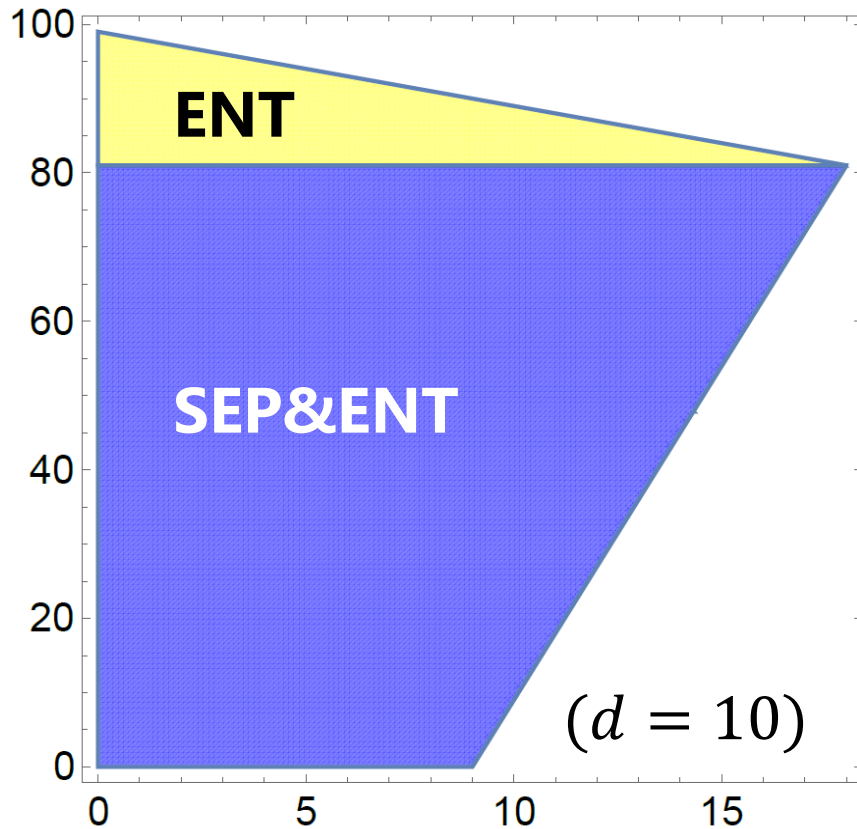
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# Geometry of two-qudit states

$$\mathcal{R}_{AB}^{(2)} = \sum \langle \lambda_i \otimes \lambda_j \rangle_{\rho_{AB}}^2 \quad \lambda_i: \text{Gell-Mann matrices}$$



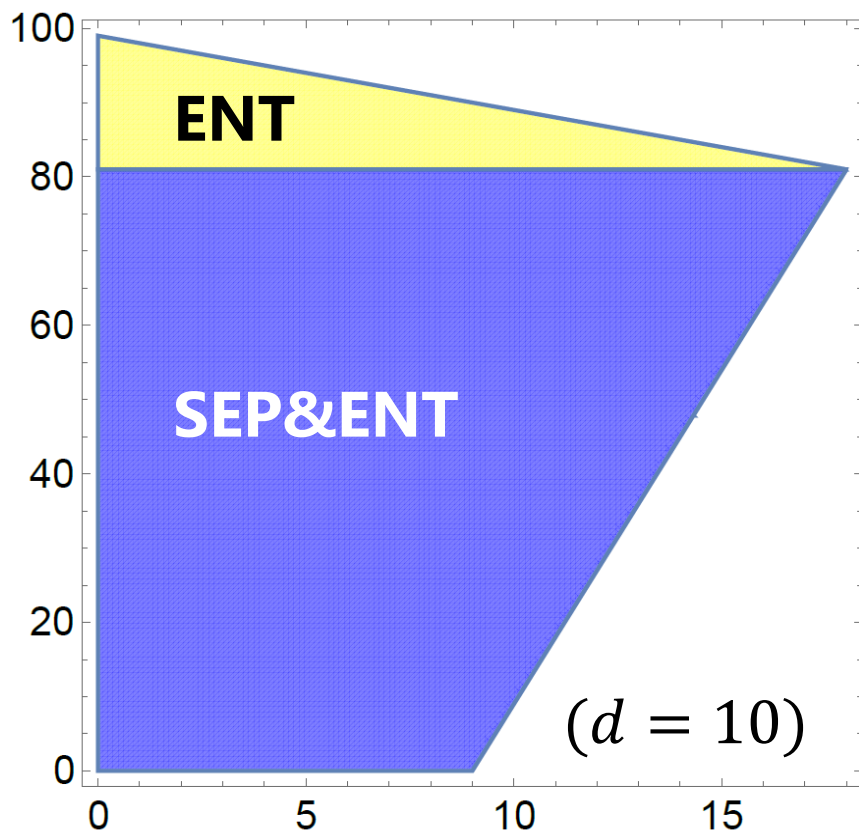
$$\mathcal{R}_{AB}^{(2)} > (d - 1)^2 \implies \text{ENT}$$

MC Tran et al, PRA 2016

$$\mathcal{R}_A^{(2)} + \mathcal{R}_B^{(2)} = \sum \langle \lambda_i \rangle_{\rho_A}^2 + \sum \langle \lambda_i \rangle_{\rho_B}^2$$

# Geometry of two-qudit states

$$\mathcal{R}_{AB}^{(2)} = \sum \langle \lambda_i \otimes \lambda_j \rangle_{\rho_{AB}}^2 \quad \lambda_i: \text{Gell-Mann matrices}$$



$$\mathcal{R}_{AB}^{(2)} > (d - 1)^2 \Rightarrow \text{ENT}$$

MC Tran et al, PRA 2016

**Questions:** Can we detect

1. entanglement optimally?
2. Schmidt number?
3. NPT entanglement?
4. bound entanglement?

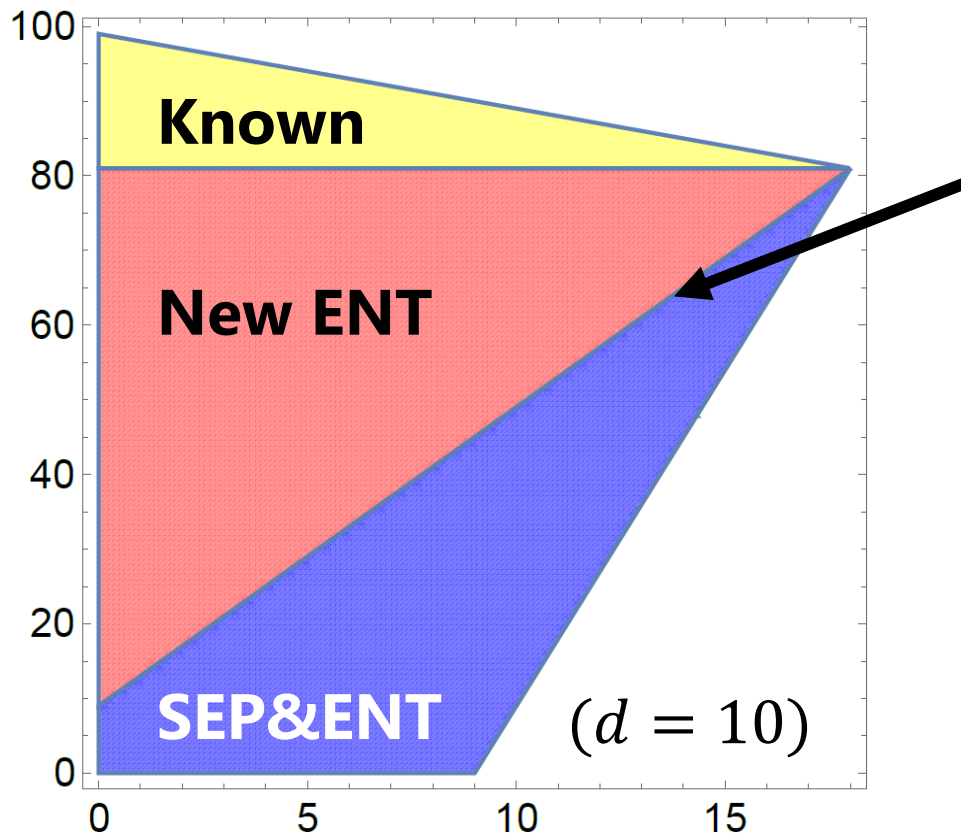
→ **We solved all of them!**

$$\mathcal{R}_A^{(2)} + \mathcal{R}_B^{(2)} = \sum \langle \lambda_i \rangle_{\rho_A}^2 + \sum \langle \lambda_i \rangle_{\rho_B}^2$$



# Result 1: Optimal separability criterion

$$\mathcal{R}_{AB}^{(2)} = \sum \langle \lambda_i \otimes \lambda_j \rangle_{\rho_{AB}}^2$$



**Result 1: Optimal line,  
which is equivalent to**

$$\text{tr}[\rho_{AB}^2] \leq \text{tr}[\rho_A^2]$$

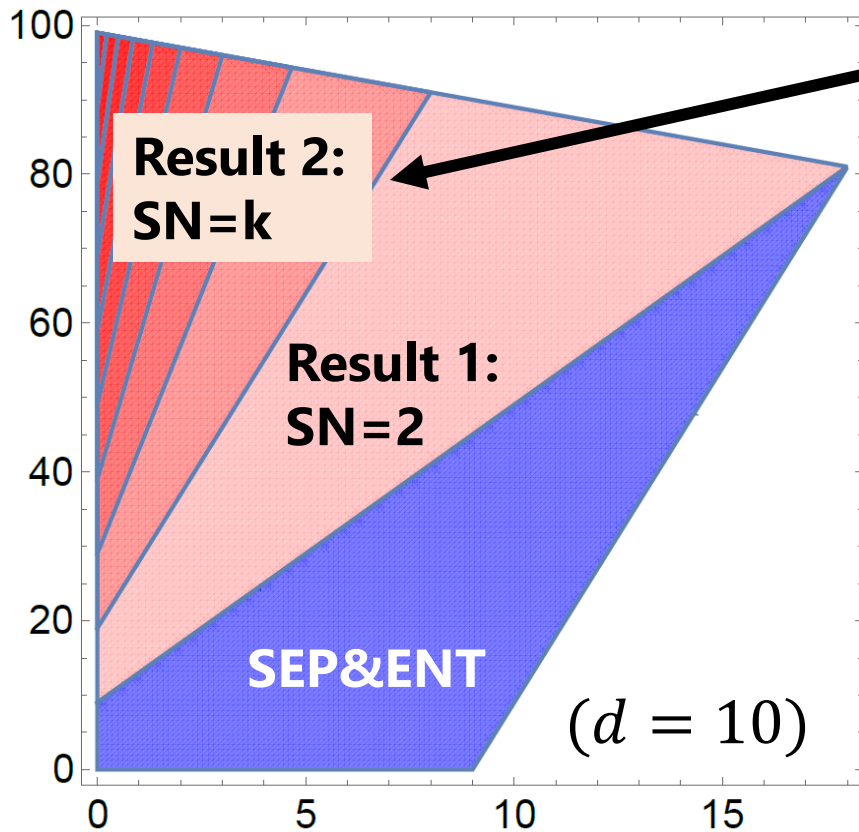
R&P&M Horodecki, PLA 1996

SI, N Wyderka, A Ketterer, O Gühne,  
PRL 2021

$$\mathcal{R}_A^{(2)} + \mathcal{R}_B^{(2)} = \sum \langle \lambda_i \rangle_{\rho_A}^2 + \sum \langle \lambda_i \rangle_{\rho_B}^2$$

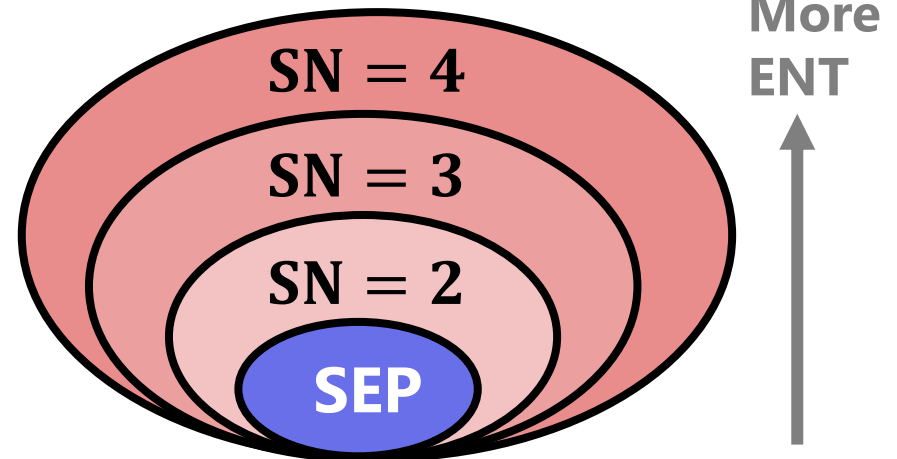
# Result 2: Schmidt number detection

$$\mathcal{R}_{AB}^{(2)} = \sum \langle \lambda_i \otimes \lambda_j \rangle_{\rho_{AB}}^2$$



*Result 2: SN lines,  
which are equivalent to*

$$\text{tr}[\rho_{AB}^2] \leq k \text{tr}[\rho_A^2]$$



$$\mathcal{R}_A^{(2)} + \mathcal{R}_B^{(2)} = \sum \langle \lambda_i \rangle_{\rho_A}^2 + \sum \langle \lambda_i \rangle_{\rho_B}^2$$

Sl, O Gühne, S Nimmrichter  
arXiv:2205.08447 (2022)

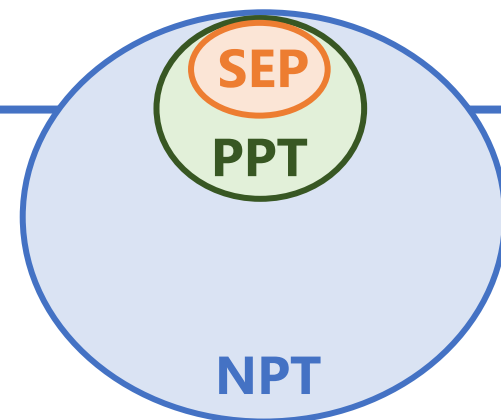
# Result 3: Efficient NPT detection

## ■ Positive Partial Transpose (PPT) criterion:

$$\rho_{AB} \in \text{SEP} \implies \rho_{AB}^{T_A} \geq 0 \text{ (called PPT state)}$$

A Peres PRL 1996; M&P&R Horodecki PLA 1996

- **Lesson:** PPT is outer approximation of SEP



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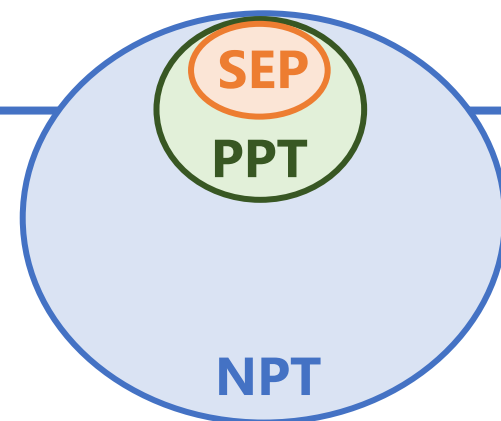
A Peres PRL 1996; M&P&R Horodecki PLA 1996

- **Lesson:** PPT is outer approximation of SEP

## ■ PT moments: $p_k = \text{tr}[(\rho_{AB}^{T_A})^k]$

- Known criterion: A Elben et al, PRL 2020

$$\rho_{AB} \in \text{PPT} \Rightarrow p_3 - p_2^2 \geq 0$$



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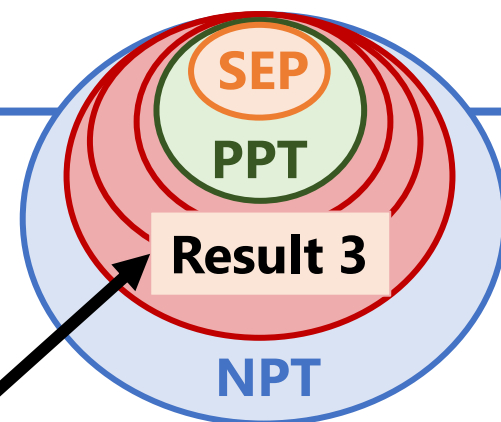
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## **Result 3: Systematic methods to detect NPT efficiently**

- **Idea:** Turn separability problem to positivity problem

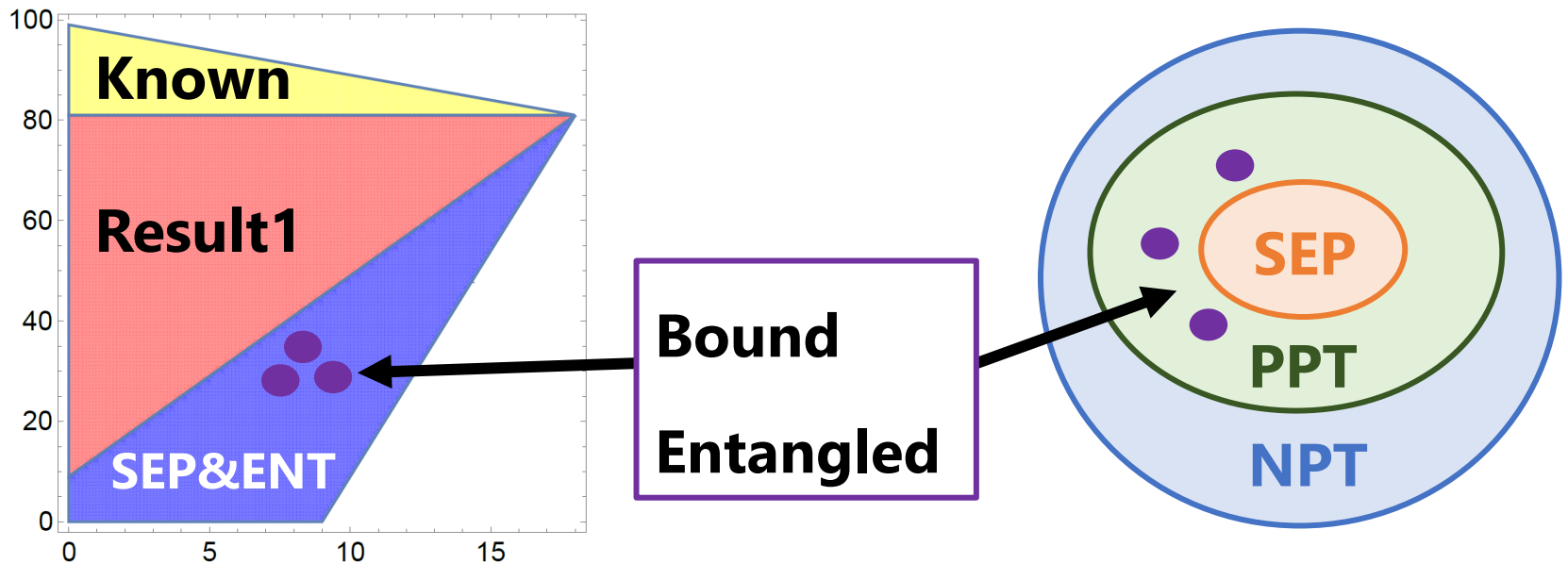
XD Yu, SL, O Gühne, PRL 2021; A Neven, et al npj QI 2021

# Bound entanglement

- Operationally it is not distillable:  $\rho_{AB} \otimes \cdots \otimes \rho_{AB} \not\Rightarrow |\text{Bell}\rangle$ 
  - ▶ very weak form of entanglement in high dimensions

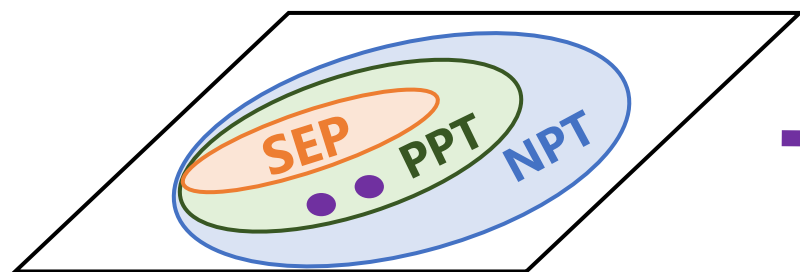
# Bound entanglement

- Operationally it is not distillable:  $\rho_{AB} \otimes \dots \otimes \rho_{AB} \not\Rightarrow |\text{Bell}\rangle$ 
  - ▶ very weak form of entanglement in high dimensions
- Mathematically it does not violate the PPT criterion
  - ▶ PPT criterion is *strictly* stronger than  $\text{tr}[\rho_{AB}^2] \leq \text{tr}[\rho_A^2]$



# Result 4: BE detection

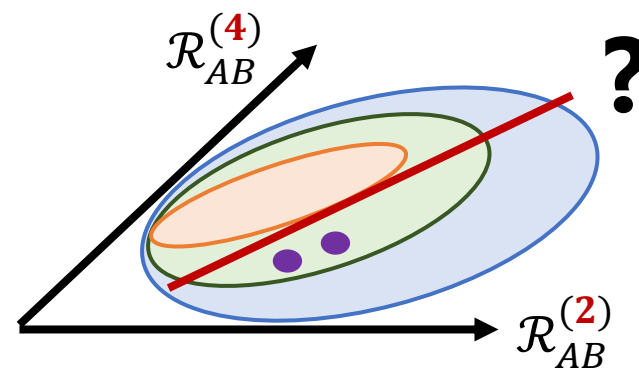
## ■ Rough idea



"positive"-separability



RMs

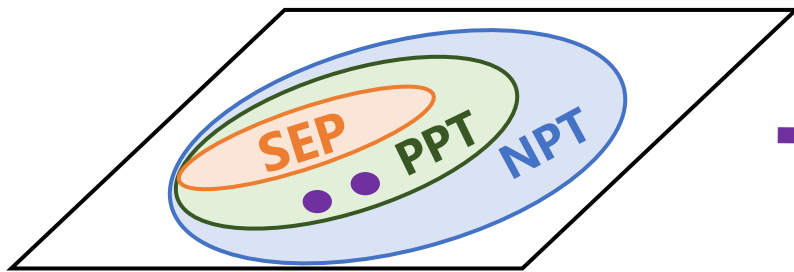


"pseudo"-separability  
de Vicente JPA 2008



# Result 4: BE detection

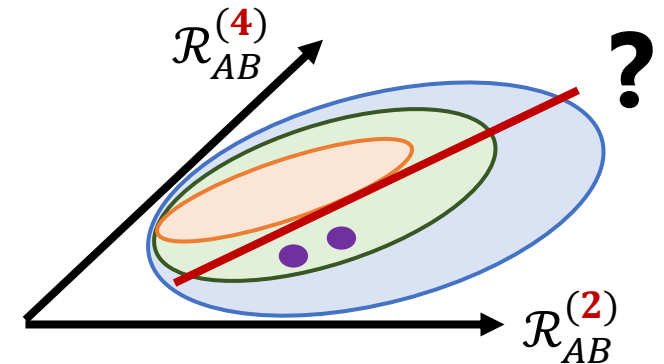
## ■ Rough idea



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de Vicente JPA 2008

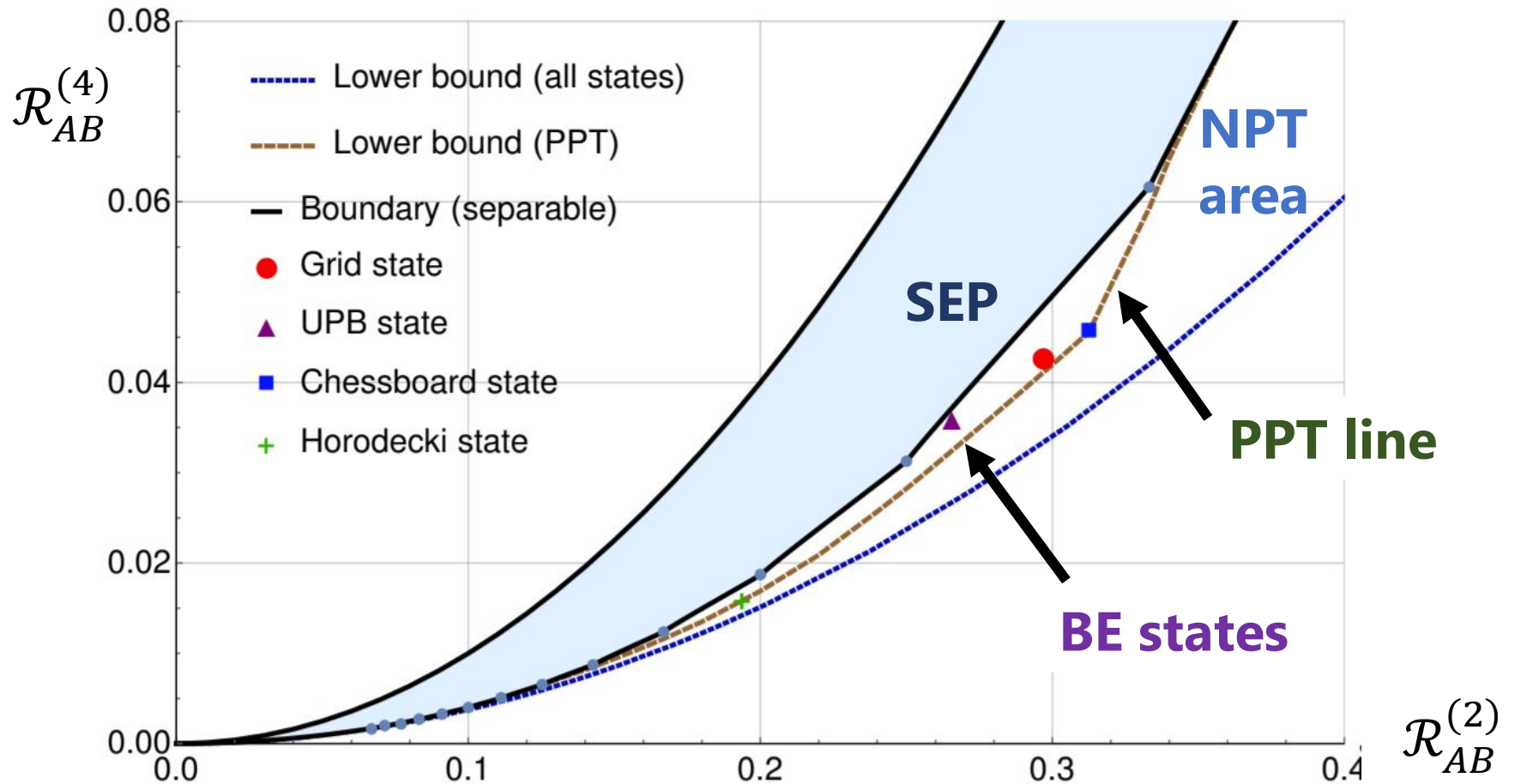
For some measurement observables, it holds that

$$\mathcal{R}_{AB}^{(r)} \stackrel{!}{=} N \int_{\mathcal{V}^{d^2-1}} d\alpha_1 \int_{\mathcal{V}^{d^2-1}} d\alpha_2 \{ \text{tr}[\varrho_{AB}(\alpha_1 \cdot \lambda_1) \otimes (\alpha_2 \cdot \lambda_2)] \}^r$$

vector of Gell-Mann matrices  
 $\downarrow$   
 $\uparrow$   
 uniformly chosen from pseudo Bloch sphere  $\mathcal{V}^{d^2-1}$

# Result 4: BE detection

## ■ Example: $3 \otimes 3$ systems



# So far so good! What's next?

✓ 1. Optimal entanglement detection

✓ 2. Schmidt number detection

✓ 3. NPT entanglement detection

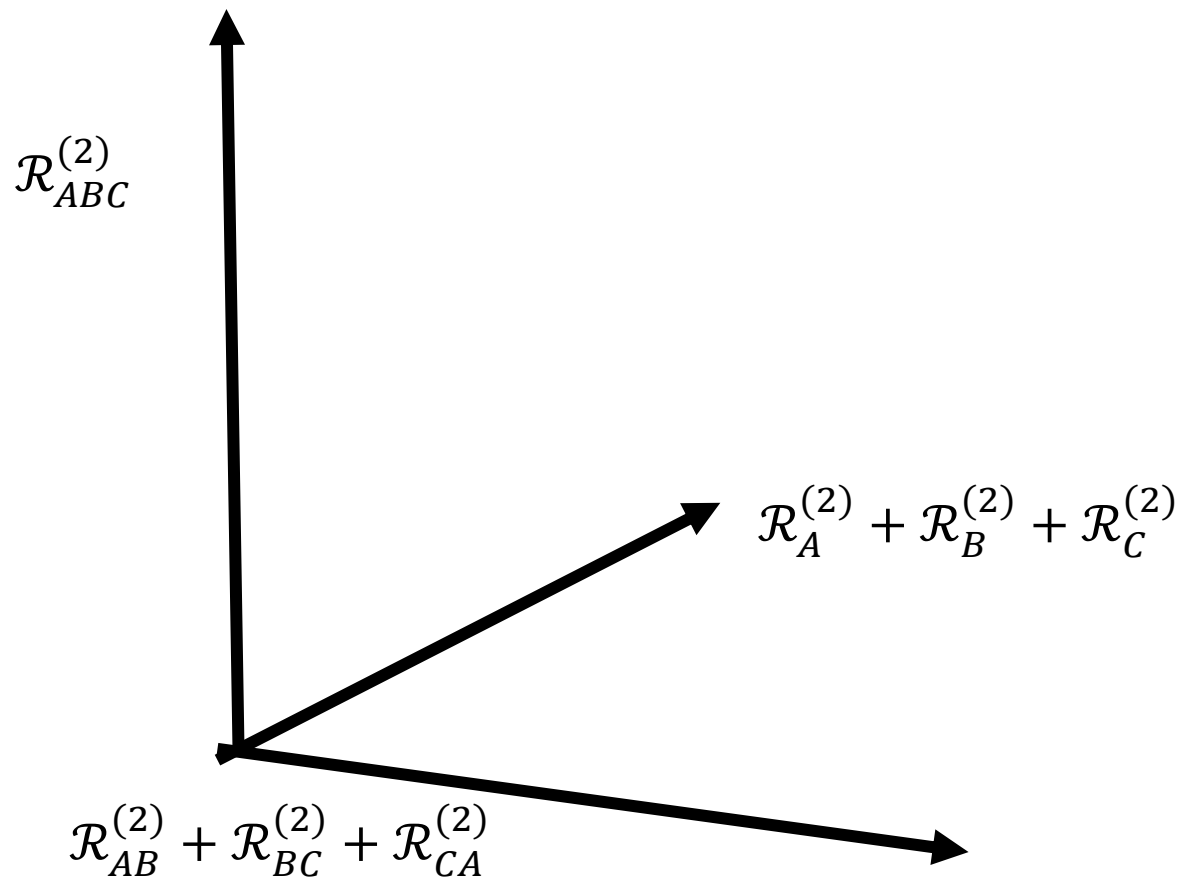
✓ 4. Bound entanglement detection

5. Three-qubit extension?

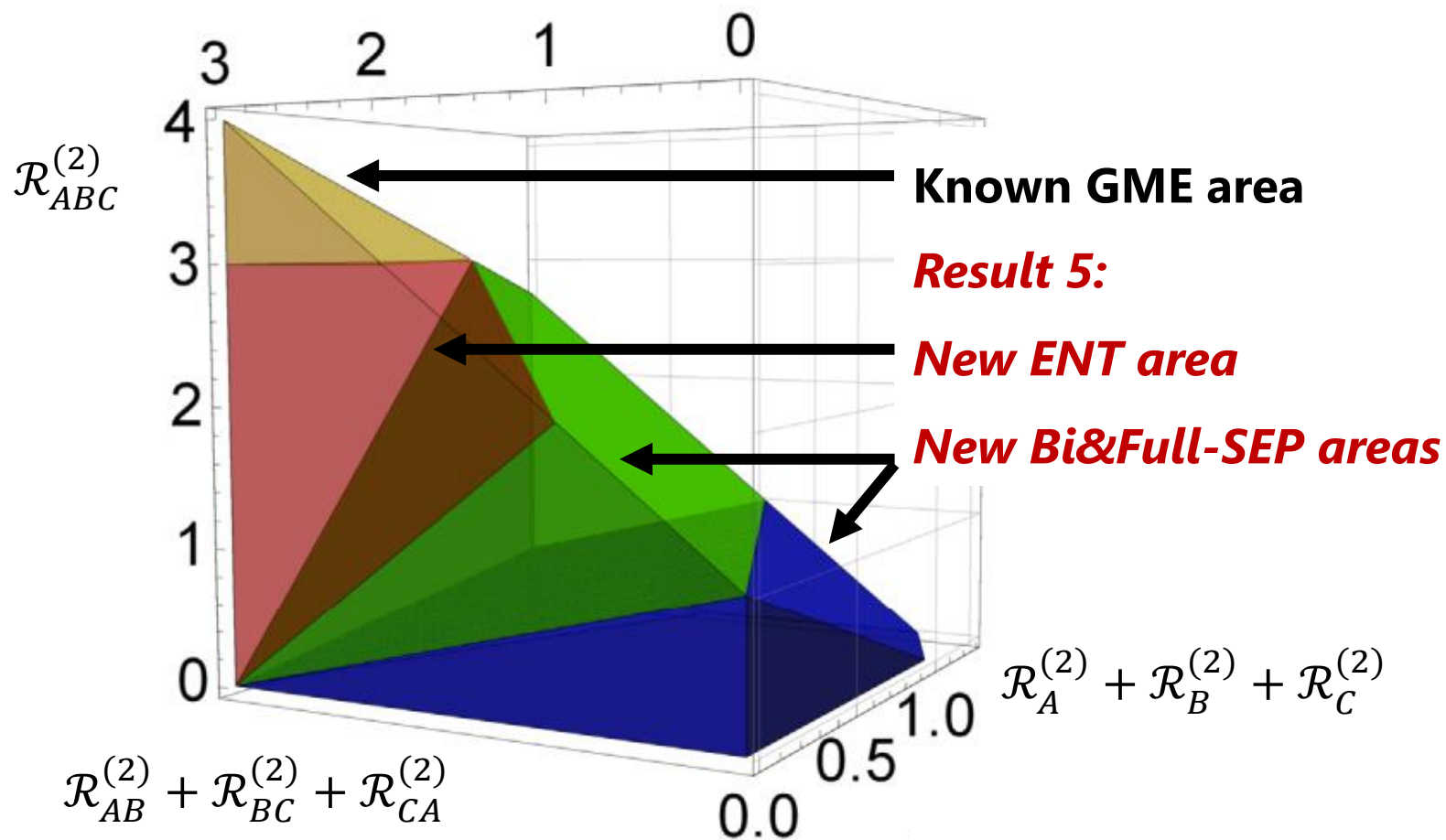
6. N-qubit extension?

7. Complete characterization?

# Result 5: Three-qubit extension



# Result 5: Three-qubit extension



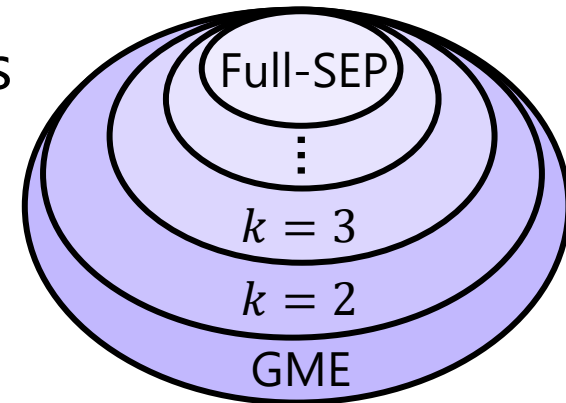
N Wyderka and O Gühne, JPA 2020;

SI, N Wyderka, A Ketterer, O Gühne, PRL 2021

# Result 6: N-qubit extension

- Consider number of separable partitions

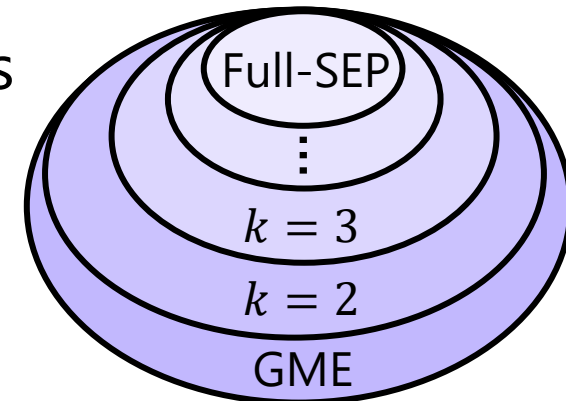
$$|\psi_{k\text{-sep}}\rangle = |\phi_1\rangle \otimes \cdots \otimes |\phi_k\rangle$$



# Result 6: N-qubit extension

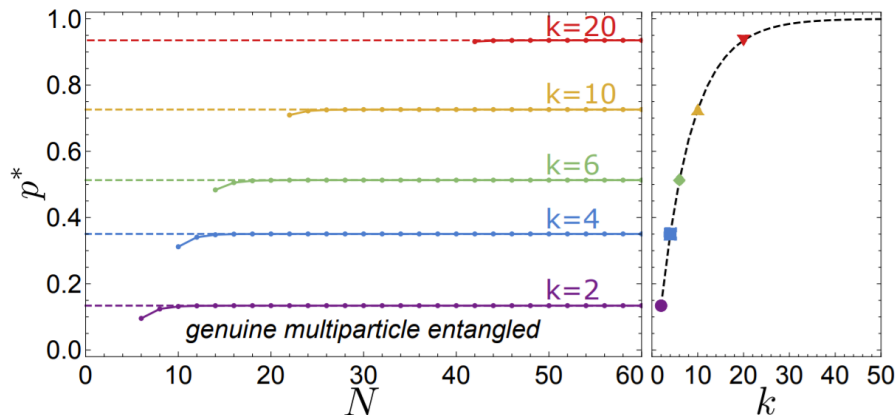
- Consider number of separable partitions

$$|\psi_{k\text{-sep}}\rangle = |\phi_1\rangle \otimes \cdots \otimes |\phi_k\rangle$$



- Result 6: Detection of  $k$ -separability**

$$\mathcal{R}_{A_1 A_2 \dots A_N}^{(2)} > f(N, k + 1) \Rightarrow \rho \text{ is } k\text{-separable}$$

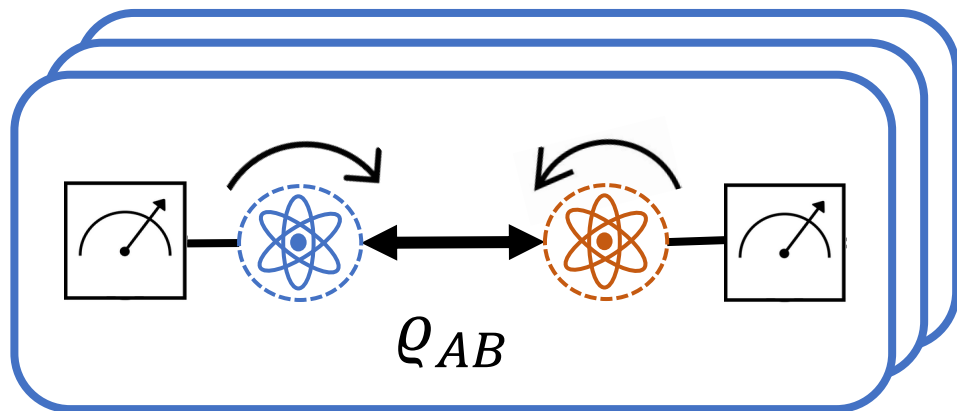


$$\rho(p) = \frac{p}{2^n} I + (1 - p) |\text{GHZ}_N\rangle\langle\text{GHZ}_N|$$

A Ketterer, [SI](#), N Wyderka, O Gühne,  
PRA (L), 2022

# Result 7: Complete characterization

## Randomized Measurements



So far

$$\alpha^2, \beta^2, \text{tr}(TT^T)$$



*Entanglement detection*

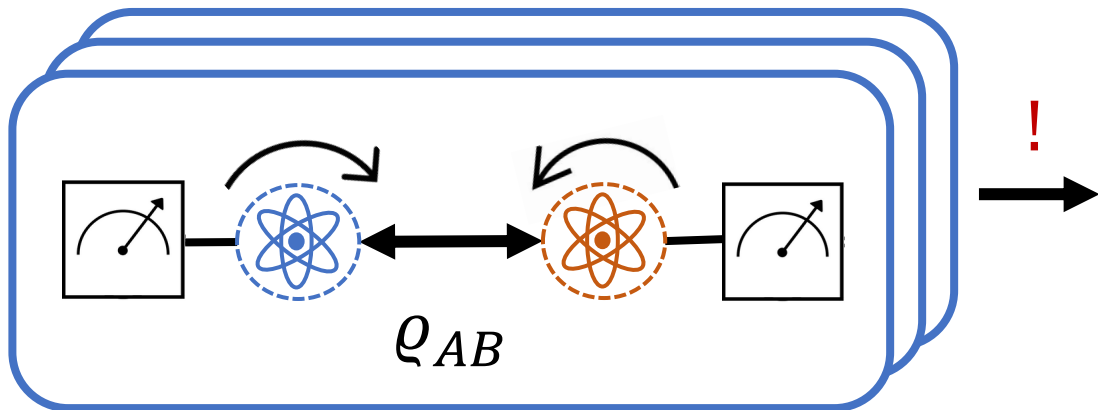
Bloch representation of two-qubit state

$$\rho_{AB} = \frac{1}{4} \left( I^{\otimes 2} + \sum \vec{\alpha} \cdot \vec{\sigma} \otimes I + \sum I \otimes \vec{\beta} \cdot \vec{\sigma} + \sum T_{ij} \sigma_i \otimes \sigma_j \right)$$



# Result 7: Complete characterization

## Randomized Measurements



## Result 7

$\alpha^2, \beta^2, \text{tr}(TT^T)$   
 $\vec{\alpha}T\vec{\beta}^T, \det(T)$   
 $[\vec{\alpha}T]^2, [\vec{\beta}T]^2, \text{tr}(TT^T TT^T),$   
*and so on...*



**teleportation fidelity,  
CHSH ineq. & PPT criterion**

coming soon...

Bloch representation of two-qubit state

$$\rho_{AB} = \frac{1}{4} \left( I^{\otimes 2} + \sum \vec{\alpha} \cdot \vec{\sigma} \otimes I + \sum I \otimes \vec{\beta} \cdot \vec{\sigma} + \sum T_{ij} \sigma_i \otimes \sigma_j \right)$$

## Summary

1. Optimal entanglement detection
2. Schmidt number detection
3. NPT entanglement detection
4. Bound entanglement detection
5. Three-qubit extension
6. N-qubit extension
7. Complete characterization

## In preparation...

1. Multiparticle BE detection?
2. Only marginals?
3. Spin squeezing?
4. Quantum metrology?

## Other my works



1. ENT change under classicalization
2. Work fluctuations and ENT
3. Q. speed limit for perturbation

**Any comments/suggestions are welcome!**

# Details of BE detection

$$1. \varrho_{AB} = \frac{1}{d^2} \left( I^{\otimes 2} + \sum \vec{a} \cdot \vec{\lambda} \otimes I + \sum I \otimes \vec{b} \cdot \vec{\lambda} + \sum T_{ij} \lambda_i \otimes \lambda_j \right)$$

2.  $\mathcal{R}_{AB}^{(4)}$  &  $\mathcal{R}_{AB}^{(2)}$  are (4<sup>th</sup> and 2<sup>nd</sup> order) polynomial functions of  $T_{ij}$

3.  $\mathcal{R}_{AB}^{(4)}$  &  $\mathcal{R}_{AB}^{(2)}$  are invariant under orthogonal rotations  $O_A, O_B$

4.  $\varrho_{AB} \in \text{SEP} \rightarrow \|T\|_{\text{tr}} \leq d - 1$  and  $\|T\|_{\text{tr}} = \|O_A T O_B^T\|_{\text{tr}} = \sum \tau_i$   
de Vicente JPA 2008

$$5. \max/\min \quad \mathcal{R}_{AB}^{(4)} = 2\sum \tau_i^4 + \left( \mathcal{R}_{AB}^{(2)} \right)^2$$

$$\text{s.t.} \quad \mathcal{R}_{AB}^{(2)} = \sum \tau_i^2$$

$$\sum \tau_i \leq d - 1$$

6. Get piecewise SEP bounds for any  $d!$

