

# Multipartite entanglement and high precision metrology

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## 1 Motivation

- Why the connection between multipartite entanglement and Fisher information is important?

## 2 Metrology and multipartite entanglement

- Quantum Fisher information
- Properties of the Quantum Fisher information
- Quantum Fisher information and entanglement

# Why the connection between multipartite entanglement and Fisher information is important?

- Genuine multipartite entanglement appears often in quantum information.
- While bipartite entanglement is quite well understood, the role of multipartite entanglement is not so clear.
- Thus, it is very interesting if we can show that it has a central role in metrology.

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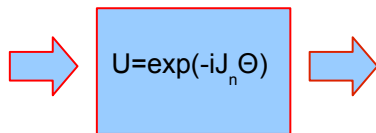
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# Metrology and multipartite entanglement in the literature

- One of the important applications of entangled multipartite quantum states is sub-shotnoise metrology.  
[V. Giovannetti, S. Lloyd, and L. Maccone, Science 306, 1330 \(2004\).](#)
- Multipartite entanglement, not simple nonseparability, is needed for extreme spin squeezing, which can be applied in spectroscopy and atomic clocks.  
[A.S. Sørensen and K. Mølmer, Phys. Rev. Lett. 86, 4431 \(2001\).](#)
- Not all entangled states are useful for phase estimation, at least in a linear interferometer.  
[P. Hyllus, O. Gühne, and A. Smerzi, 82, 012337 \(2009\).](#)

# Quantum Fisher information

- Let us consider the following process:



- The dynamics described above is  $\rho_{\text{out}} = e^{-i\theta J_{\vec{n}}} \rho e^{+i\theta J_{\vec{n}}}$ .
- We would like to determine the angle  $\theta$  by measuring  $\rho_{\text{out}}$ .

# Quantum Fisher information II

## Quantum Cramér-Rao bound

The phase estimation sensitivity is limited as

$$\Delta\theta \geq \frac{1}{\sqrt{F_Q[\varrho, J_{\vec{n}}]}},$$

where  $F_Q$  is the quantum Fisher information,  $\varrho$  is a quantum state and  $J_{\vec{n}}$  is a collective angular momentum component.

- The Braunstein-Caves quantum Fisher information is

$$F[\varrho, X] = \sum_{ij} \frac{2(\lambda_i - \lambda_j)^2}{\lambda_i + \lambda_j} |X_{ij}|^2.$$

C.W. Helstrom, *Quantum Detection and Estimation Theory* (1976),

A. S. Holevo, *Probabilistic and Statistical Aspect of Quantum Theory* (1982).

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# Properties of the Quantum Fisher information

Two important properties:

- 1 For a pure state  $\rho$ , we have  $F[\rho, J_I] = 4(\Delta J_I)_\rho^2$ .
- 2  $F[\rho, J_I]$  is convex in the state, that is  
 $F[p_1\rho_1 + p_2\rho_2, J_I] \leq p_1F[\rho_1, J_I] + p_2F[\rho_2, J_I]$ .

It also follows that  $F[\rho, J_I] \leq 4(\Delta J_I)_\rho^2$ .

C.W. Helstrom, *Quantum Detection and Estimation Theory* (1976).

A. S. Holevo, *Probabilistic and Statistical Aspect of Quantum Theory* (1982).

S.L. Braunstein and C.M. Caves, *Phys. Rev. Lett.* 72, 3439 (1994).

L. Pezzé and A. Smerzi, *Phys. Rev. Lett.* 102, 100401 (2009).

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# Quantum Fisher information and entanglement

Pezzé, Smerzi, PRL 2009

For  $N$ -qubit separable states we have

$$F_Q[\varrho, J_I] \leq N.$$

Here,  $J_I = \frac{1}{2} \sum_{k=1}^N \sigma_I^{(k)}$  where  $\sigma_I^{(k)}$  are the Pauli spin matrices. The maximum for the left-hand side is  $N^2$ .

Thus, **for separable states**

$$\Delta\theta \geq \frac{1}{\sqrt{N}},$$

while **for entangled states**

$$\Delta\theta \geq \frac{1}{N}.$$

## Observation 1

For  $N$ -qubit separable states we have

$$\sum_{l=x,y,z} F_Q[\varrho, J_l] \leq 2N. \quad (1)$$

- Eq. (1) is a condition for the average sensitivity of the interferometer. All states violating Eq. (1) are entangled.

GT, PRA 85, 022322 (2012); P. Hyllus et al., PRA 85, 022321 (2012).

## Observation 2

For quantum states we have the bound

$$\sum_{l=x,y,z} F_Q[\rho, J_l] \leq N(N+2). \quad (2)$$

GHZ states and  $N$ -qubit symmetric Dicke states with  $\frac{N}{2}$  excitations saturate Eq. (2).

- Dicke states have been investigated recently in several experiments.
- In general, pure symmetric states for which  $\langle J_l \rangle = 0$  for  $l = x, y, z$  saturate Eq. (2).

# Quantum Fisher information and multipartite entanglement

Next, we will consider  $k$ -producible or  $k$ -entangled states:

## Observation 3

For  $N$ -qubit  $k$ -producible states

$$\sum_{l=x,y,z} F_Q[\rho, J_l] \leq nk(k+2) + (N-nk)(N-nk+2).$$

where  $n$  is the integer part of  $\frac{N}{k}$ . For the  $k = N - 1$  case, this bound can be improved

$$\sum_{l=x,y,z} F_Q[\rho, J_l] \leq N^2 + 1. \quad (3)$$

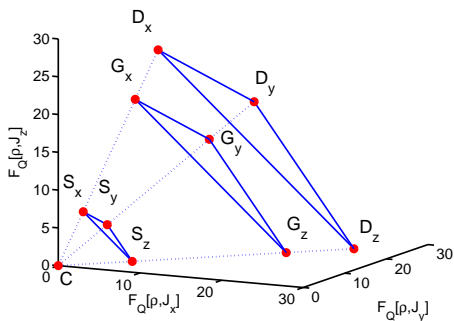
Eq. (3) is also the inequality for biseparable states. **Any state that violates Eq. (3) is genuine multipartite entangled.**

# Quantum Fisher information and multipartite entanglement II

## Fact

*Genuine multipartite entanglement, not simple nonseparability is needed to achieve maximum sensitivity in a linear interferometer.*

# Quantum Fisher information and multipartite entanglement III



**Figure:** Points in the  $(F_Q[\rho, J_x], F_Q[\rho, J_y], F_Q[\rho, J_z])$ -space for  $N = 6$ .

- Points corresponding to separable states are not above the  $S_x - S_y - S_z$  plane.
- Points corresponding to biseparable states are not above the  $G_x - G_y - G_z$  plane.



# Which part of the space corresponds to quantum states? - Points

- A completely mixed state

$$\rho_C = \frac{\mathbb{1}}{2^N}.$$

corresponds to the point  $C(0, 0, 0)$ .

- States corresponding to the point  $S_x(0, N, N)$  is

$$|\Psi\rangle_{S_x} = \left| +\frac{1}{2} \right\rangle_x^{\otimes N/2} \otimes \left| -\frac{1}{2} \right\rangle_x^{\otimes N/2}.$$

$S_y$  and  $S_z$  are similar.

# Which part of the space corresponds to quantum states? - Points II

- $D_z$  :  $N$ -qubit symmetric Dicke state with  $\frac{N}{2}$  excitations.

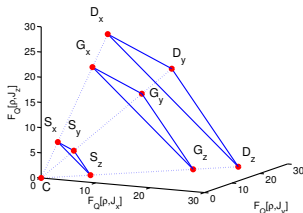
$$|\mathcal{D}_N^{(N/2)}\rangle = \binom{N}{\frac{N}{2}}^{-\frac{1}{2}} \sum_k \mathcal{P}_k \{ |0\rangle^{\otimes \frac{N}{2}} \otimes |1\rangle^{\otimes \frac{N}{2}} \},$$

where  $\sum_k \mathcal{P}_k$  denotes summation over all possible permutations.

- $N$ -qubit GHZ states

$$|\Psi\rangle_{GHZ_z} = \frac{1}{\sqrt{2}} (|0\rangle^{\otimes N} + |1\rangle^{\otimes N}).$$

# Which part of the space corresponds to quantum states? - 2D polytopes

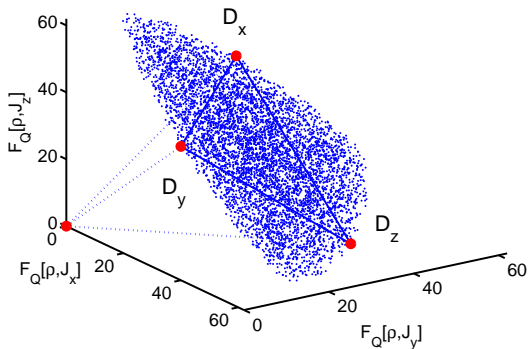


- For all points in the  $S_x, S_y, S_z$  polytope, there is a corresponding pure product state for even  $N$ .
- For given  $F[\rho, J_l]$  for  $l = x, y, z$ , such a state is defined as

$$\rho = \left[ \frac{\mathbb{1}}{2} + \frac{1}{2} \sum_{l=x,y,z} c_l \sigma_l \right]^{\otimes N/2} \otimes \left[ \frac{\mathbb{1}}{2} - \frac{1}{2} \sum_{l=x,y,z} c_l \sigma_l \right]^{\otimes N/2},$$

where  $c_l^2 = 1 - \frac{F_Q[\rho, J_l]}{N}$ , where  $\sum_l c_l^2 = 1$ .

# Which part of the space corresponds to quantum states? - 2D polytopes II



**Figure:** Randomly chosen points in the  $(F_Q[\rho, J_x], F_Q[\rho, J_y], F_Q[\rho, J_z])$ -space corresponding to states  $|\Psi(\alpha_x, \alpha_y, \alpha_z)\rangle$  for  $N = 8$ .

- All the points are in the plane of  $D_x$ ,  $D_y$  and  $D_z$ .

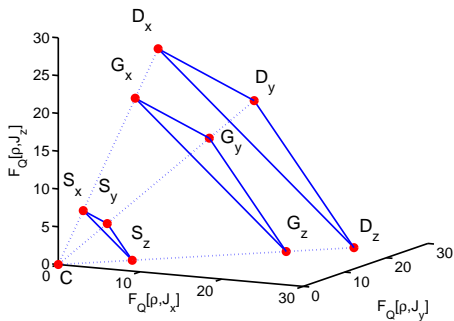
# Which part of the space corresponds to quantum states? - 3D polytopes

- A pure state mixed with the completely mixed state

$$\varrho^{(\text{mixed})}(\rho) = \rho\varrho + (1 - \rho)\frac{\mathbb{1}}{2^N}$$

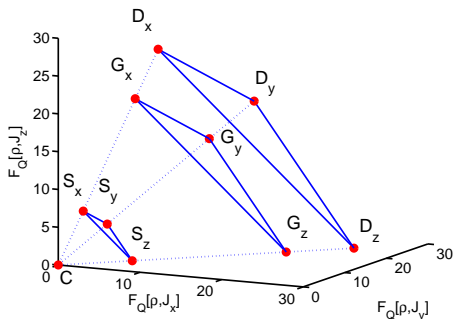
- The states  $\varrho^{(\text{mixed})}(\rho)$  are on a straight line on our figures.

# Which part of the space corresponds to quantum states? - 3D polytopes II



**Observation 5.** If  $N$  is even, then there is a separable state for each point in the  $S_x, S_y, S_z, C$  polytope.

# Which part of the space corresponds to quantum states? - 3D polytopes III



**Observation 6.** If  $N$  is divisible by 4, then for all the points of the  $D_x$ ,  $D_y$ ,  $D_z$ ,  $G_x$ ,  $G_y$ ,  $G_z$  polytope, there is a quantum state with genuine multipartite entanglement.

# Summary

- We defined entanglement conditions in terms of the quantum Fisher information.
- We showed that genuine multipartite entanglement is needed for maximum metrological sensitivity.

See:

G. Tóth, PRA 85, 022322 (2012).

Similar paper: Hyllus et al, PRA 85, 022321(2012); Krischek et al., PRL 107, 080504 (2011).

THANK YOU FOR YOUR ATTENTION!

